

# Thermal convection and the Sun's supergranulation

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# Outline

- 1 Introduction : the case of supergranulation
- 2 How to progress ?

# The Sun's supergranulation

## The Doppler image

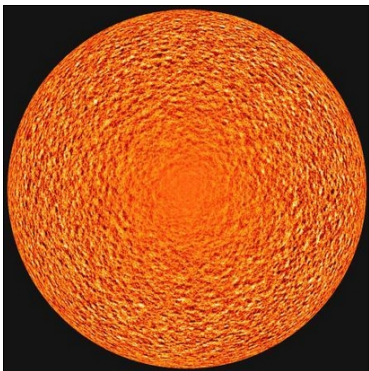


Fig.: from SOHO/MDI.

# The Sun's supergranulation

The spherical harmonic spectrum : SOHO/MDI

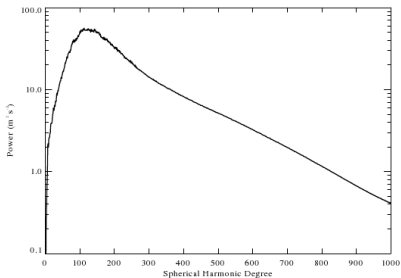
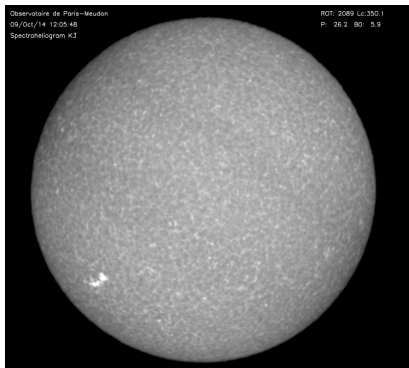


Fig.: From Hathaway et al. 2000 ;  $L_{\text{max}}=120$ ,  $\lambda_{SG} = 36.4\text{Mm}$ .

# The Sun's supergranulation

## The network in the chromosphere



**Fig.:** Image of the Sun at  $\lambda = 393.37\text{nm}$ ,  $\text{Ca}^+$  K3 line (Meudon observatory).

# The Sun's supergranulation

## The network in the photosphere

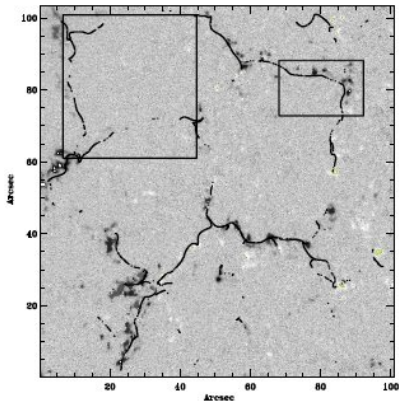


Fig.: Network magnetic fields (Roudier et al. 2009).

# Problems with supergranulation

- This scale singles out of a continuum
- No consensus about its origin or models
- Obviously some connection with magnetic fields
- Is it a universal feature of stellar convection ?

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# Various explanations

- Simon & Leighton (1964) :  $\text{He}^{2+} + \text{e}^{-} \rightarrow \text{He}^{+} + \text{heat}$
- Cloutman (1979) : Rip currents
- Rieutord et al. (2000) : Large-scale instability of surface convection favoured by strong stratification
- Rincon & Rieutord (2003) : Linear instability from fixed-flux BC, scaled limited by  $\vec{B}$ -field.

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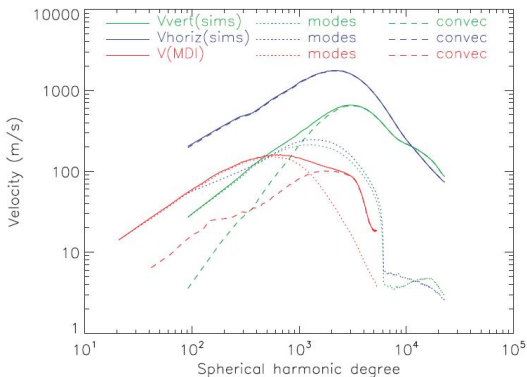
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## Three attempts by DNS

- Rieutord, Ludwig et al. (2002), resol.  $315^2 \times 96$  ( $30^2 \times 3\text{Mm}^3$ )
- Rincon, Lignières and Rieutord (2005), resol.  $1024^2 \times 82$
- Stein et al. (2008), resol.  $1000^2 \times 500$  ( $96^2 \times 20\text{Mm}^3$ )

**Large amount of computing power but no sign of supergranulation !**

# Spectra from DNS



**FIG.:** LES within a box  $48 \times 48 \times 20\text{Mm}^3$ , resol.  $500^3$  (Georgobiani et al. 2007).



# Direct modeling

## Some ideas

Granulation forms in the thermal boundary layer where the radiative cooling is strong and the entropy gradient strong as well.

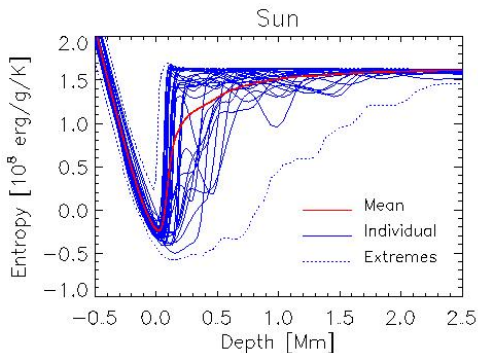


FIG.: From Bob Stein simulations.

# Direct modeling

## Some ideas (2)

- Below the thermal boundary layer, the gradient is still superadiabatic but the flux is fixed.
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# Thermal convection at fixed flux

## First results

- Hurle et al. (1966) showed that when boundaries are fully insulating, namely conductivity  $\rightarrow 0$ , then  $k_{\text{crit}} \rightarrow 0$ . Unstable scales are larger and larger horizontally.
- In fact, it turns out that  $k = 0$  is the most unstable mode. But how this result changes when a vertical magnetic field is imposed ?

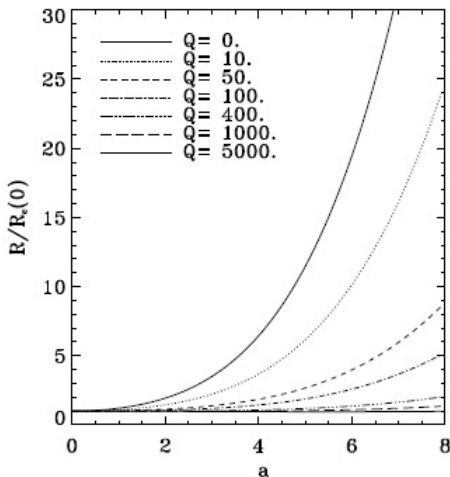
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## Fixed flux and magnetic field

This case was investigated by Rincon and Rieutord in 2003.



# Fixed flux and magnetic field

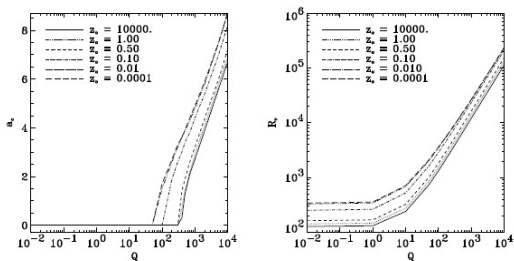


FIG.: Critical Rayleigh number and critical wavenumber as a function of the magnetic field.

$Q = (B_o^2 d^2) / (\mu_o \mu \eta)$  is the Chandrasekhar number.

# Fixed flux and magnetic field

- This model shows that **there is a critical magnetic beyond which the critical wavenumber is finite**. It is characterized by the Chandrasekhar number  $Q=61$  for a very strong stratification typical of the Sun.
- Putting solar numbers shows that the supergranulation scale is selected if the magnetic field is in the range 100G to 1kG at 5Mm deep.



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# Outlooks

The preceding numbers are in the right orders of magnitude. They suggest investigating this instability in a more realistic framework :

- Use the spherical geometry
- Introduce a slight rotation to check the drift of convective cells
- Use entropy diffusion instead of heat (temperature) diffusion.

In fact follow up the early work of Glatzmaier and Gilman 1981 and the recent work of Jones et al. 2009.

**END**