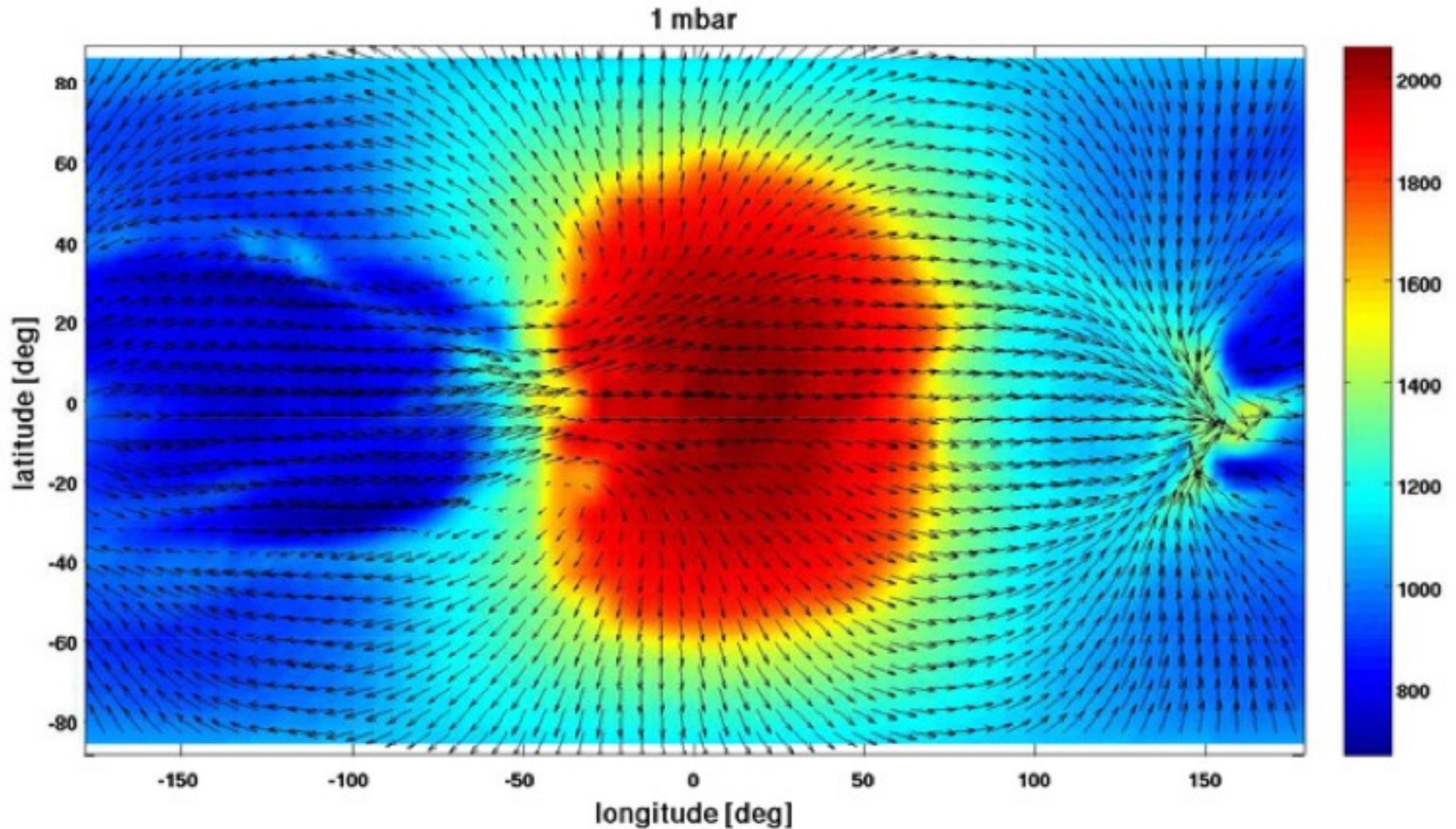
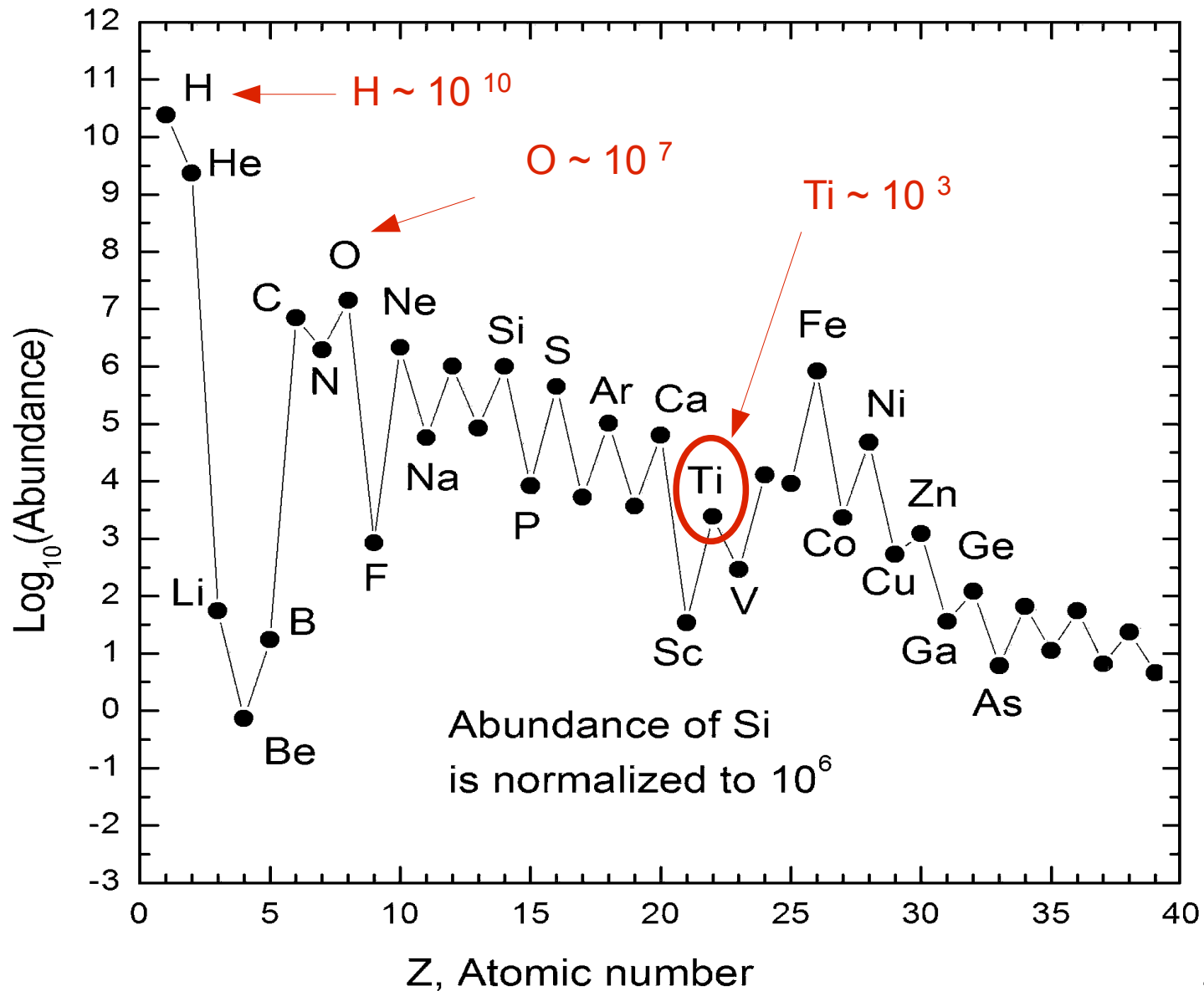


# Day/night cold trap of TiO in hot Jupiters atmospheres

*Vivien Parmentier  
with Adam Showman*



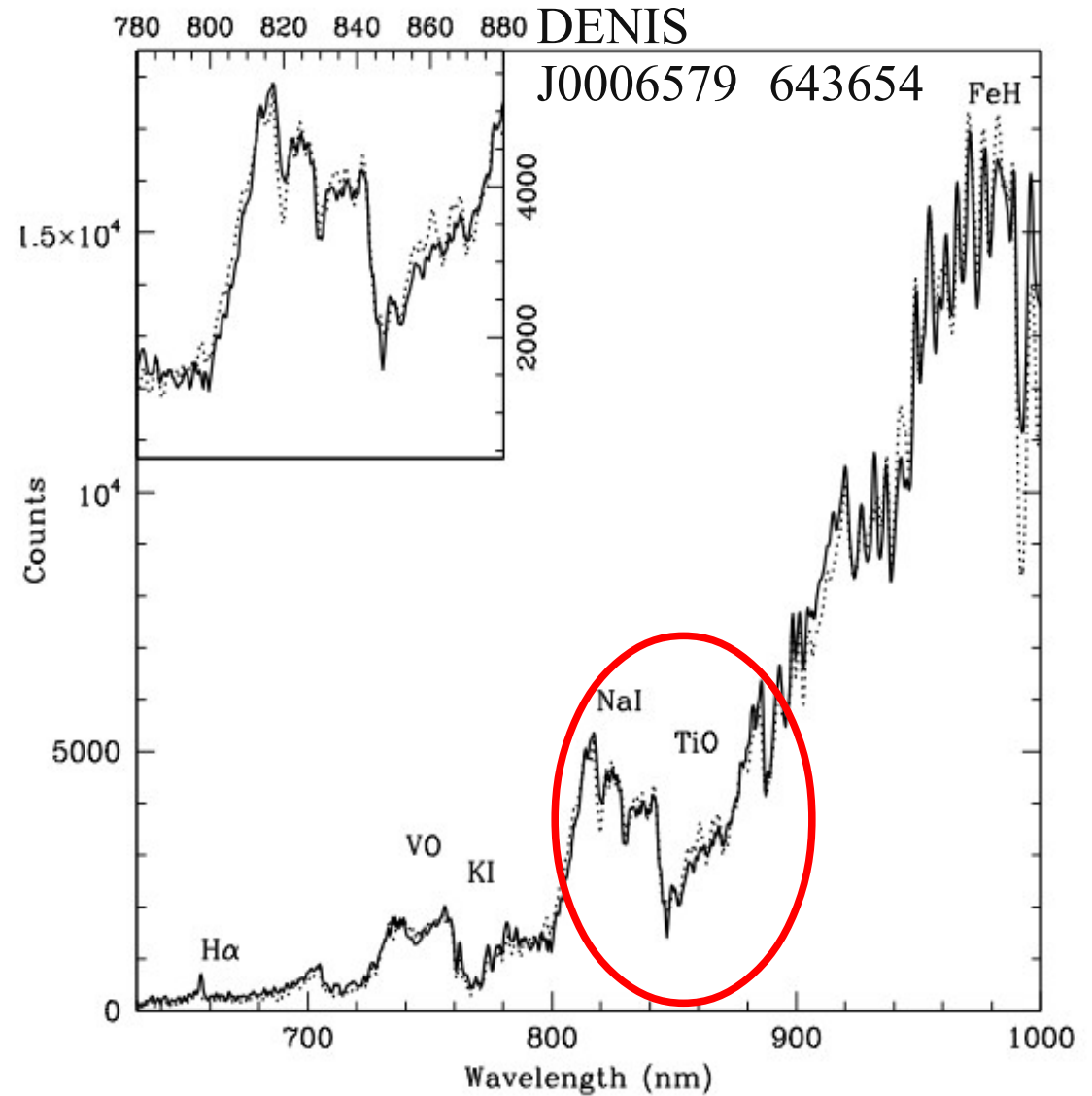
# Background : Abundance of Ti the solar nebula.



# Background : presence of TiO/VO in brown dwarf

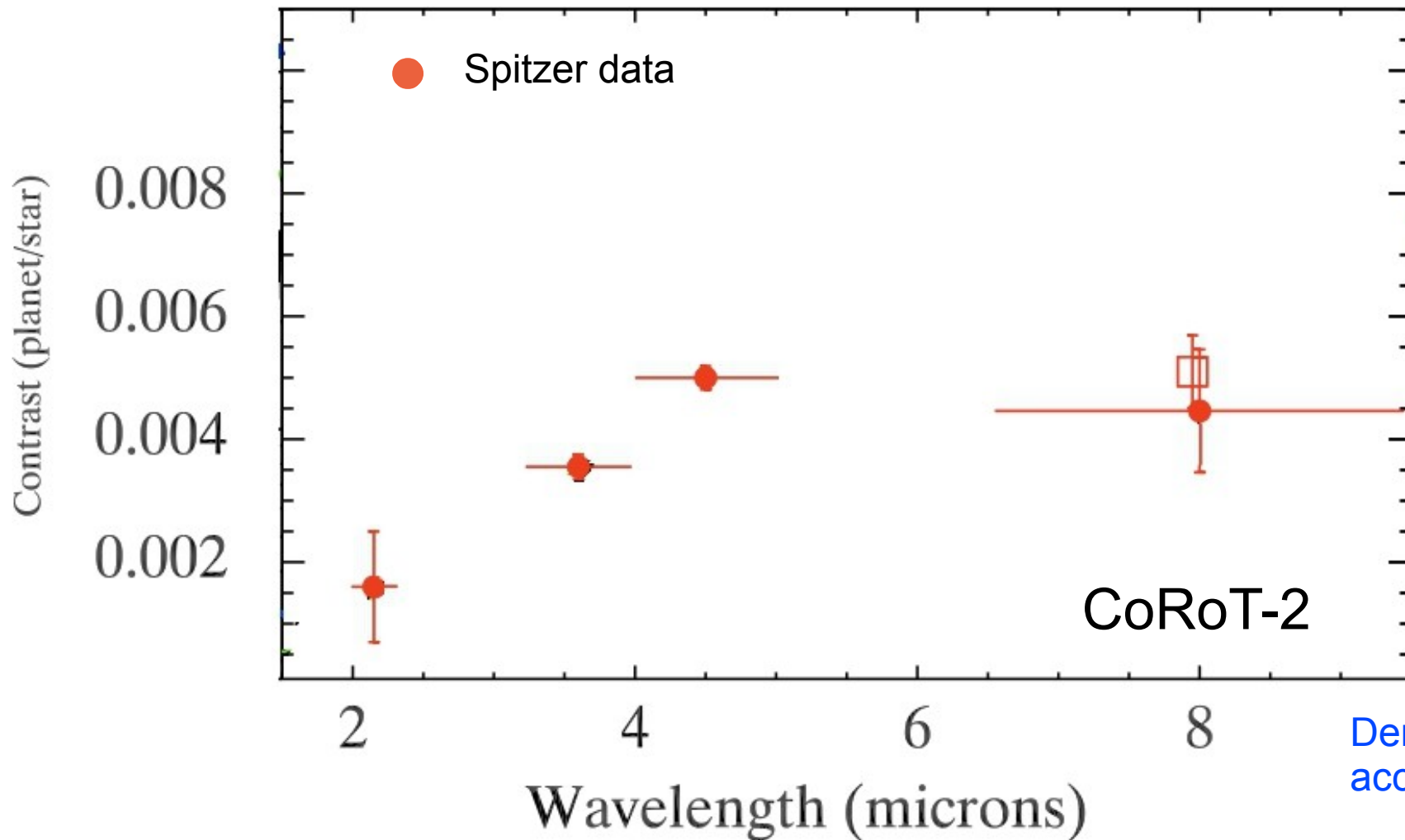
TiO : Titanium oxyde

The presence of the TiO absorption band define the separation between M dwarf and L dwarf.



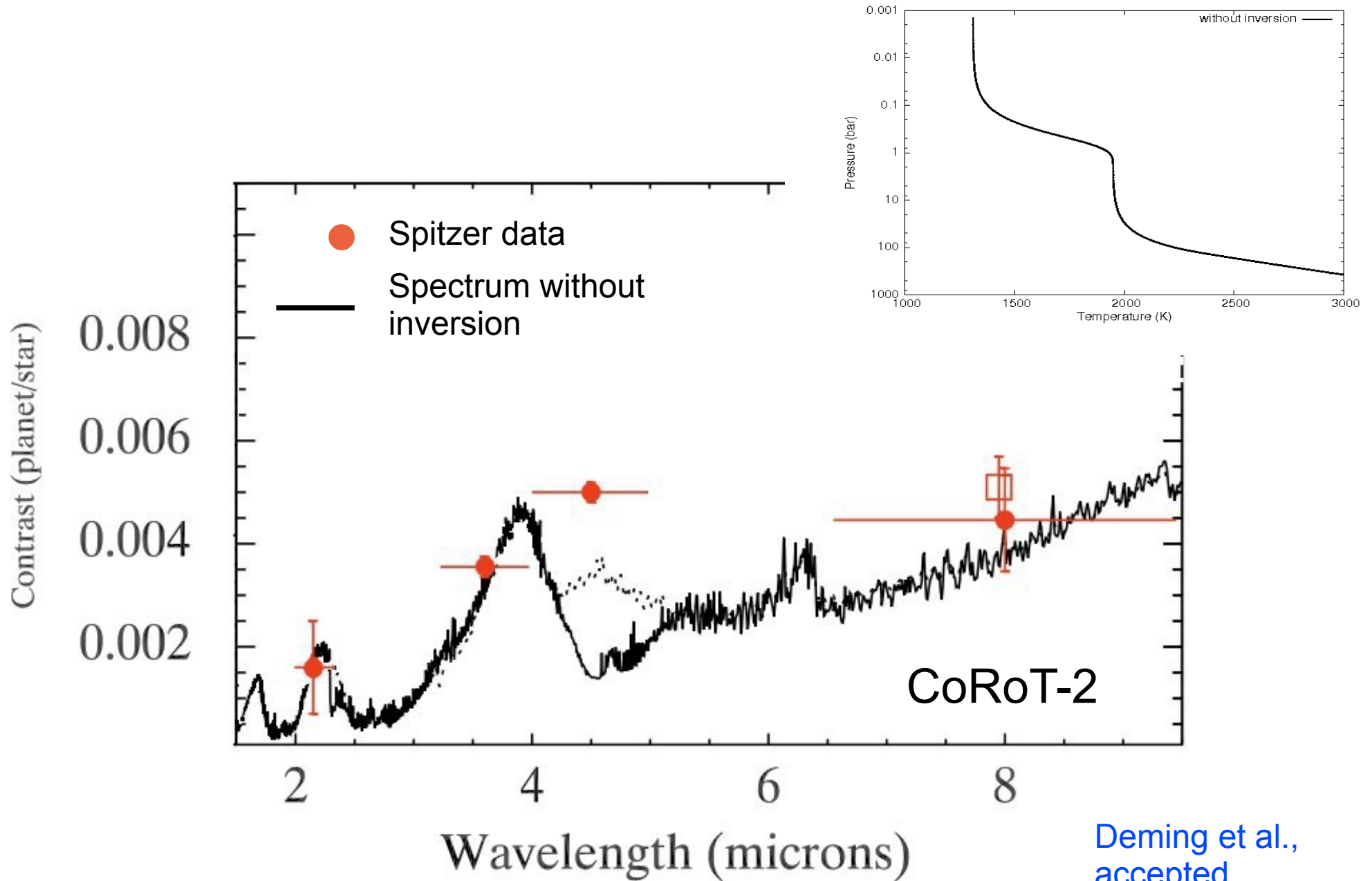
# Background : presence of TiO in hot Jupiter atmospheres ?

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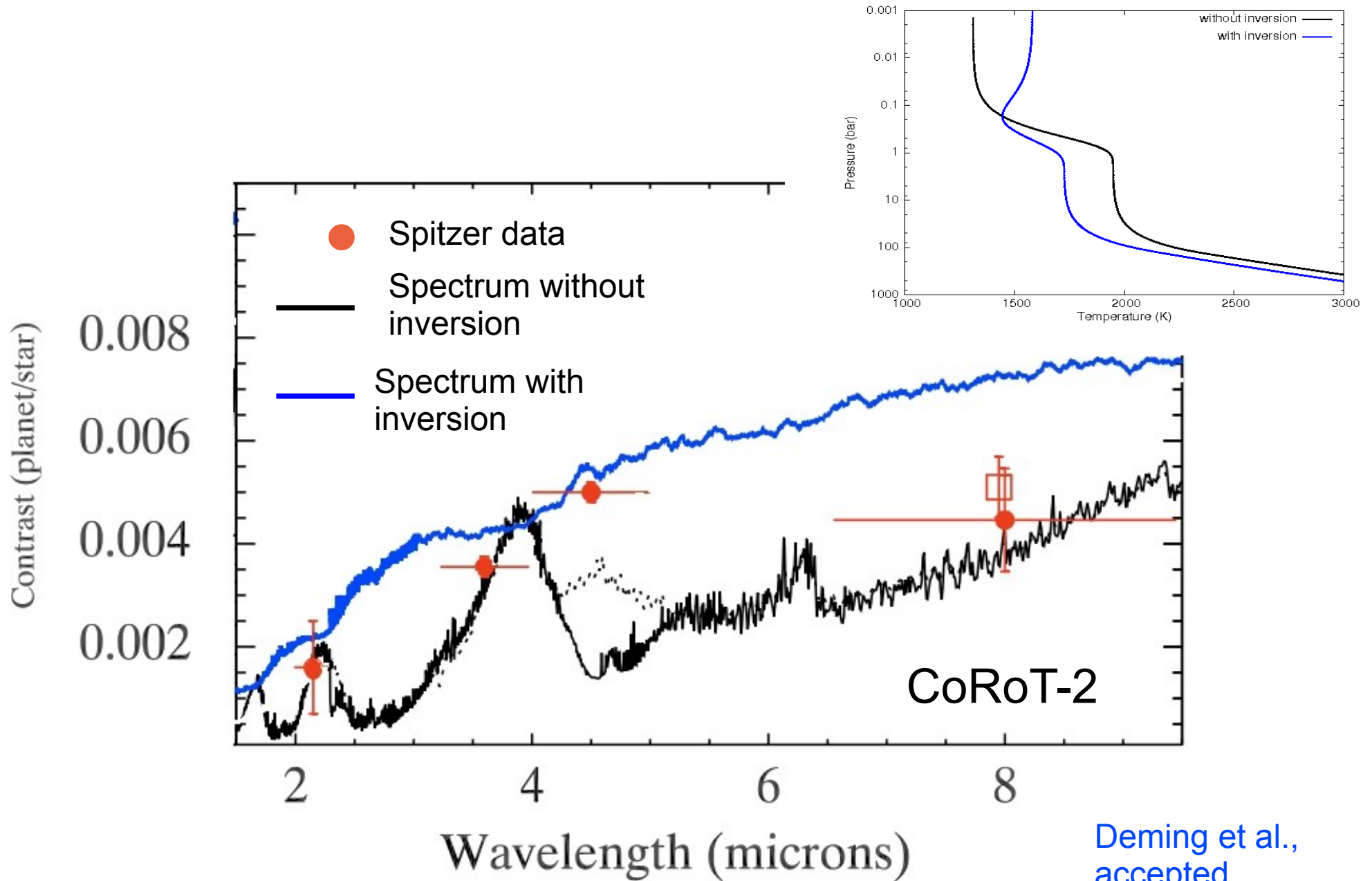
Deming et al.,  
accepted

# Background : presence of TiO in hot Jupiter atmospheres ?

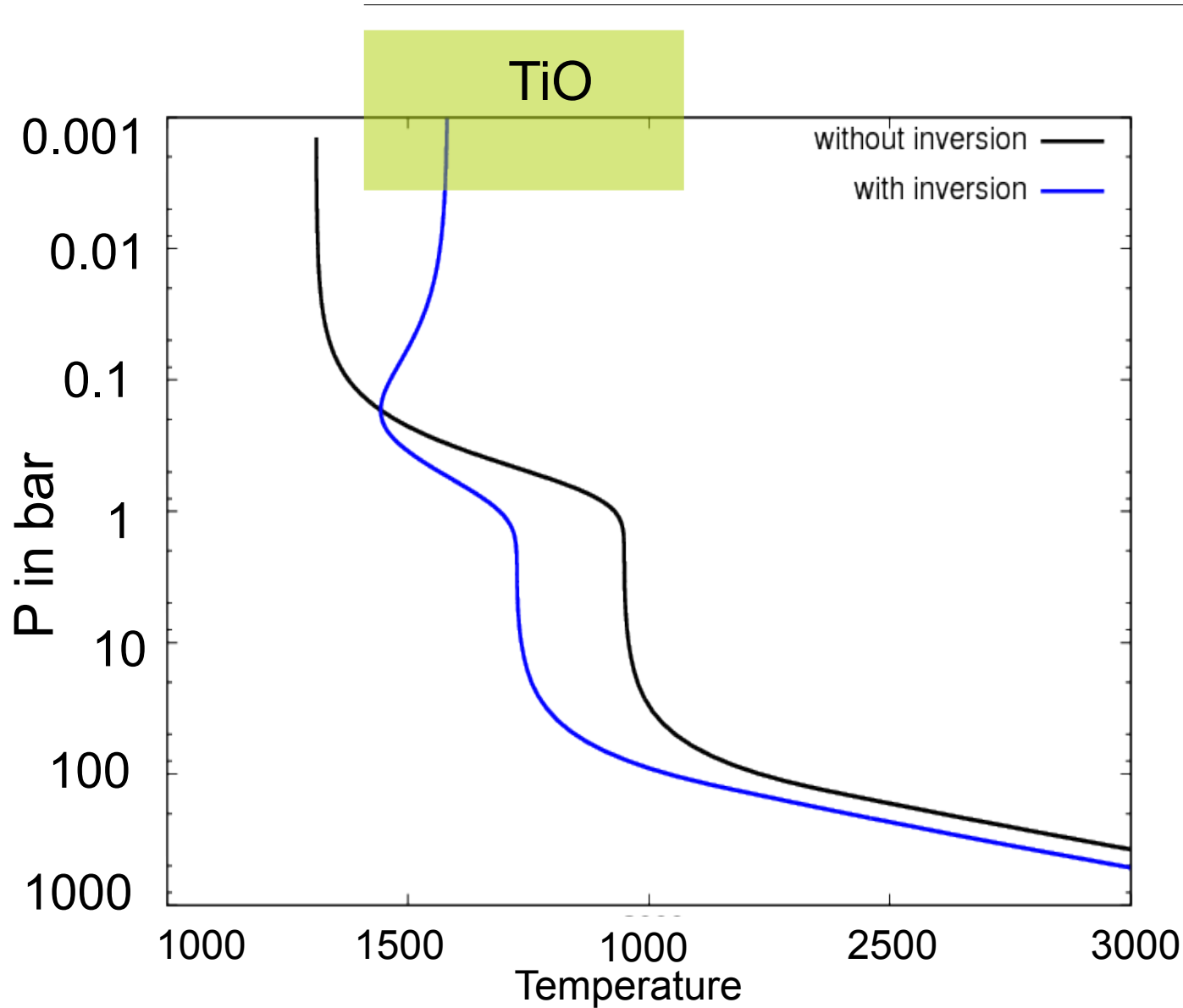


Deming et al.,  
accepted

# Background : presence of TiO in hot Jupiter atmospheres ?



# Background : Temperature inversion due to TiO



The presence of a visible absorber in the upper atmosphere warm it and create a temperature inversion

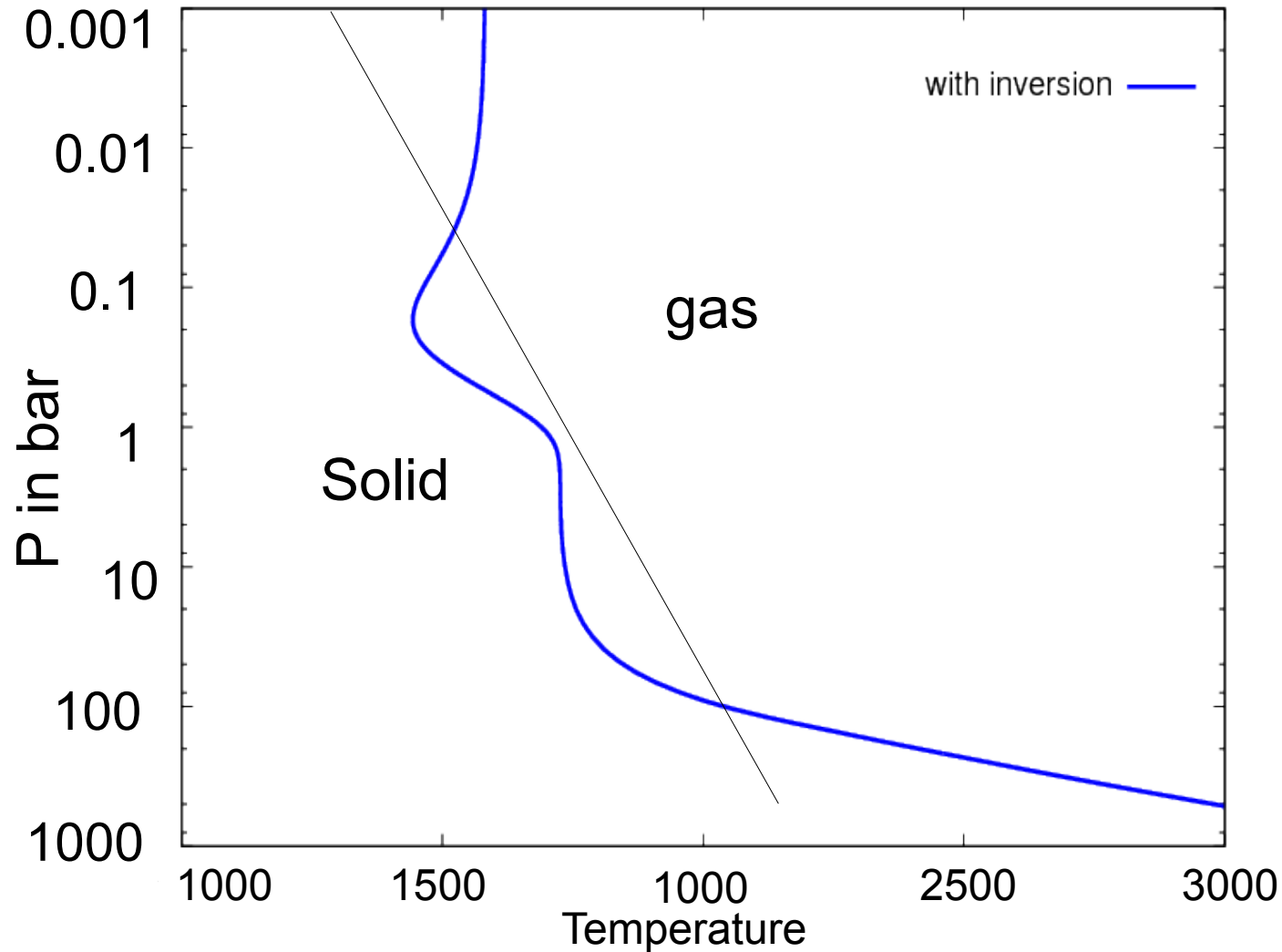
→ Condition :  $f_{\text{TiO}} = 0.5 f_0$  between 1mBar and 0.1mBar

(Spiegel et al. 2009)

# Background : Vertical cold trap

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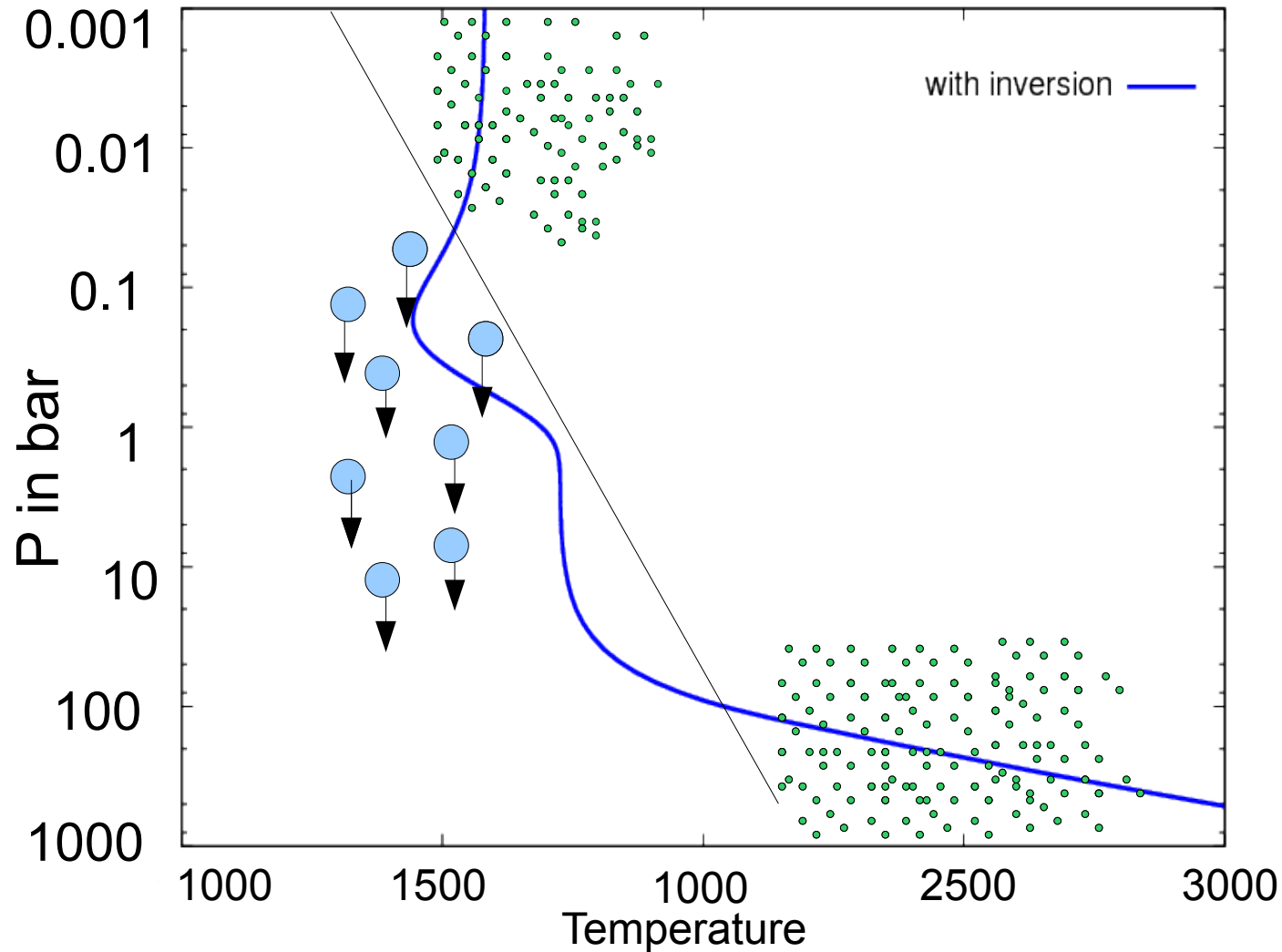
Particles diffuse to the cold region, condense and settle down.





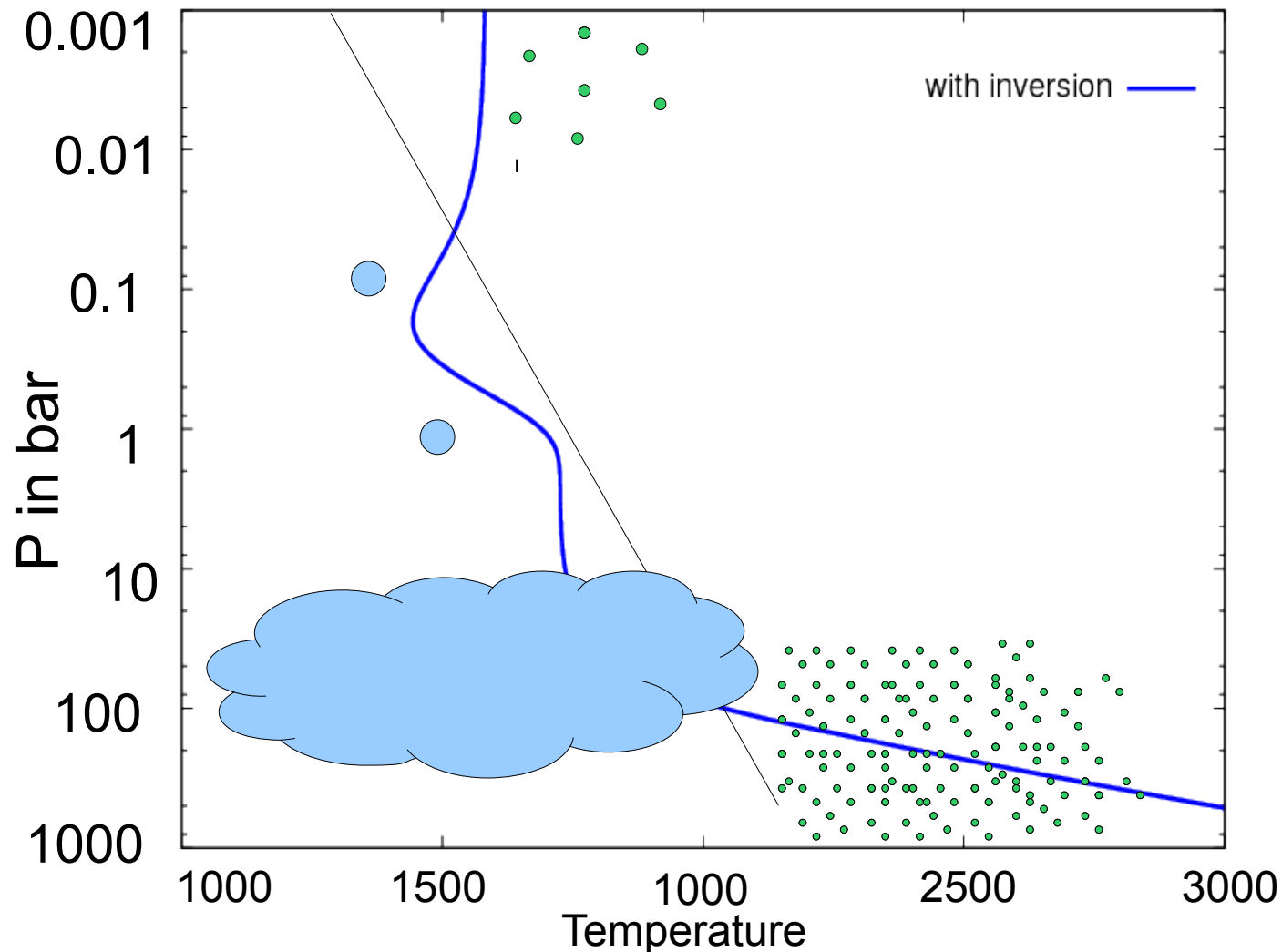
# Background : Vertical cold trap

Particles diffuse to the cold region, condense and settle down.



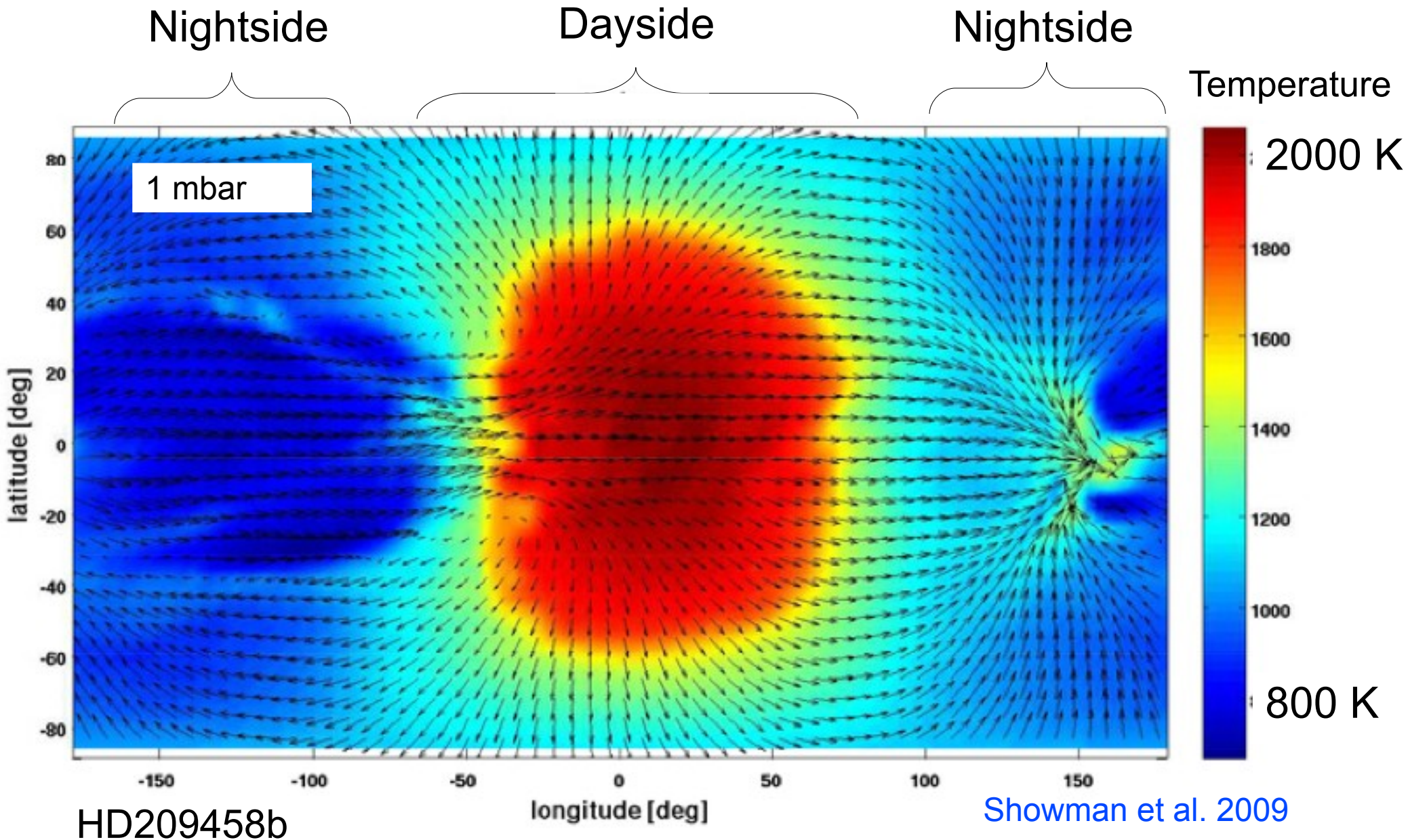
# Background : Vertical cold trap

Particles diffuse to the cold region, condense and settle down.



The high atmosphere is depleted in the species.

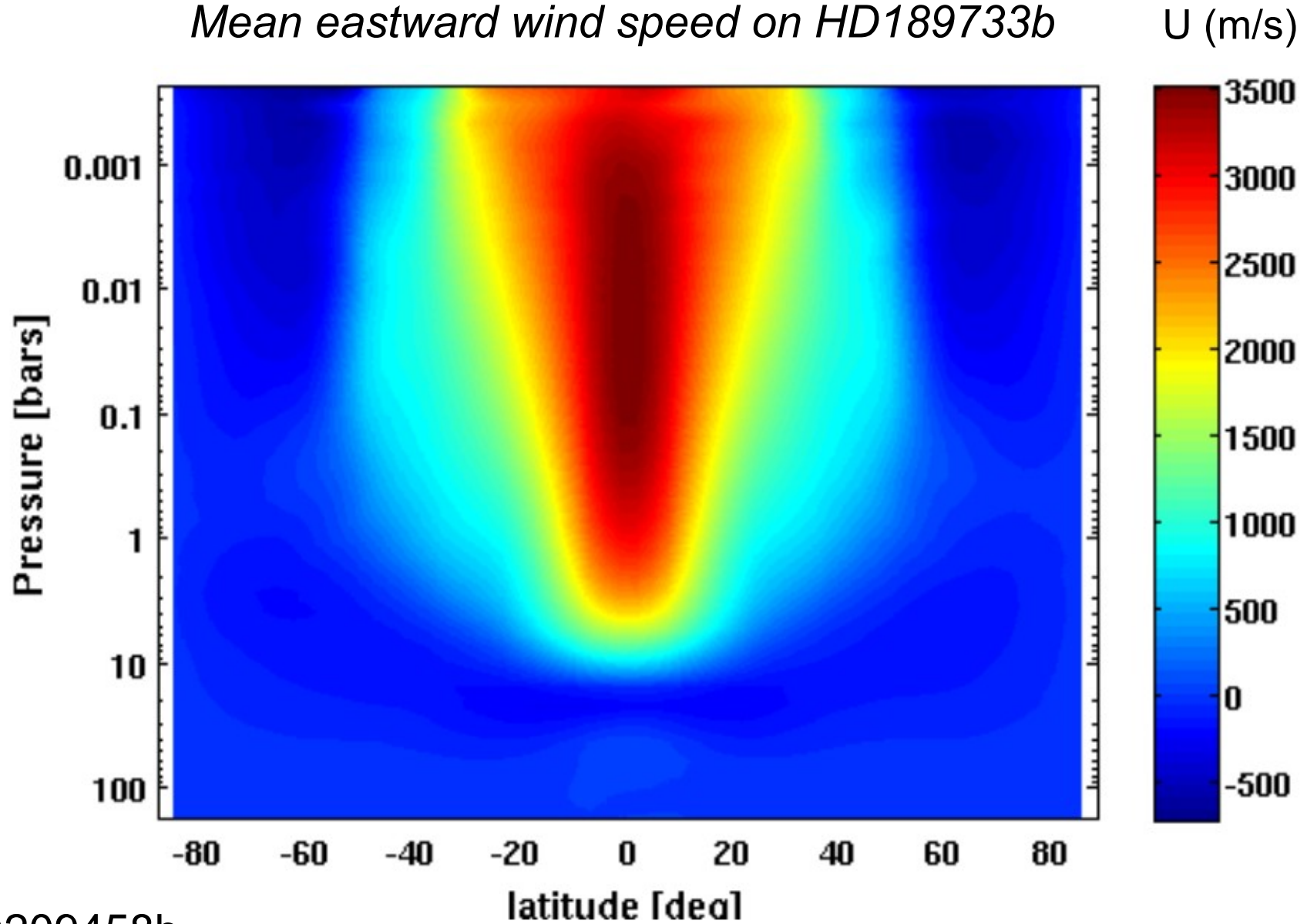
# Background : Day/night temperature contrast on hot Jupiter



# Background : Superrotation on hot Jupiter

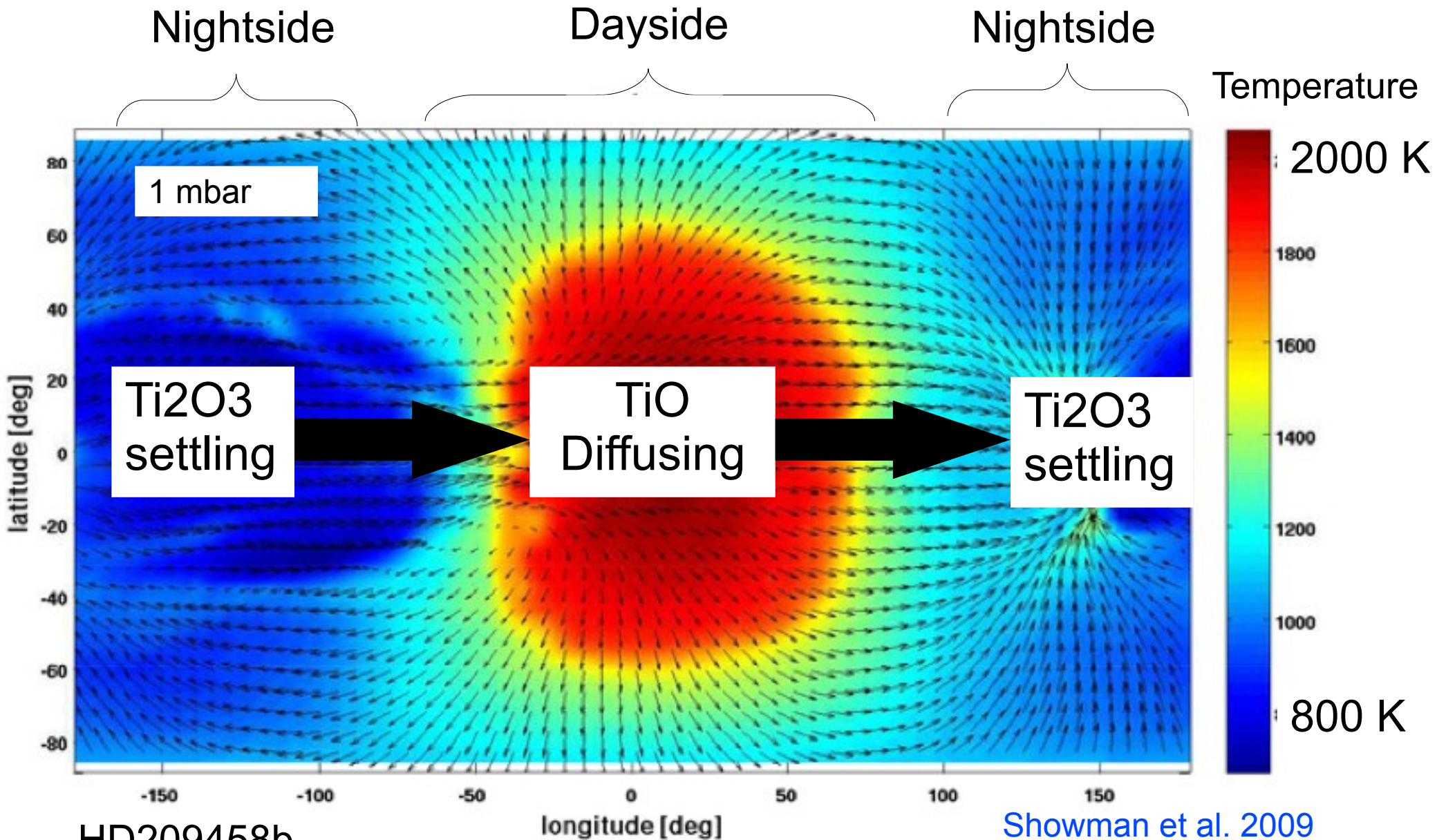
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*Mean eastward wind speed on HD189733b*

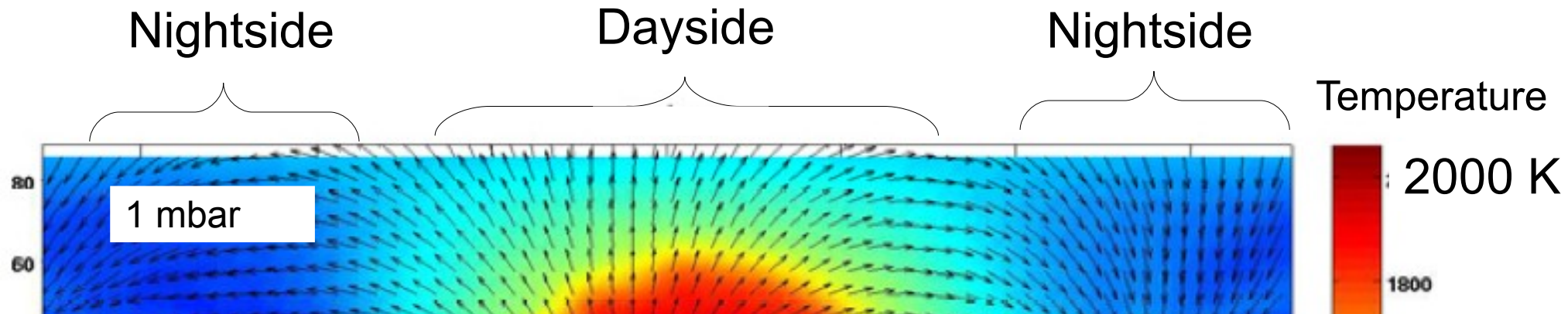


HD209458b

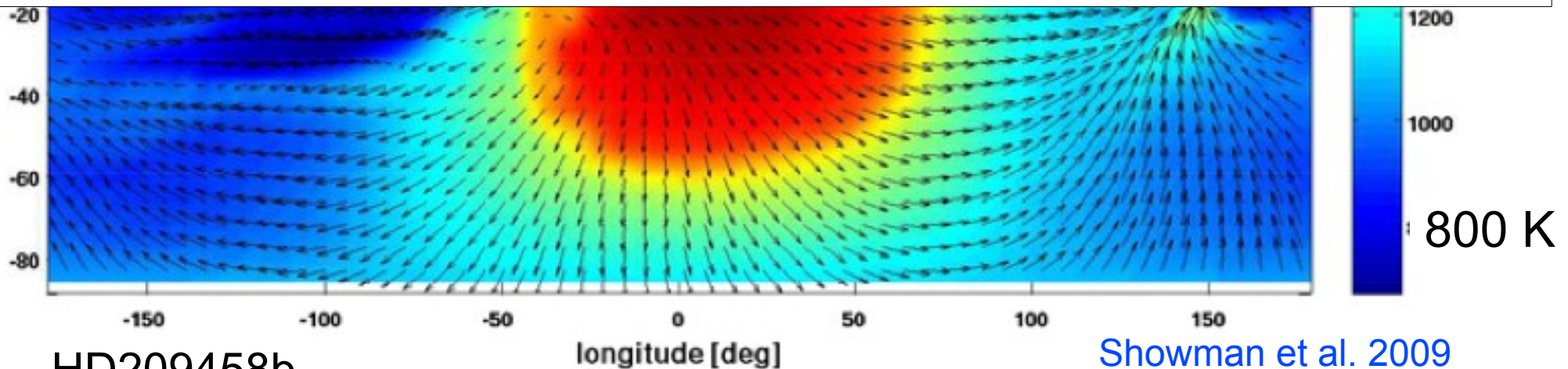
# Background : Day/night cold trap on hot Jupiter



# Background : Day/night cold trap on hot Jupiter



How the atmospheric circulation and the day/night temperature contrast will affect the equilibrium concentration of TiO ?

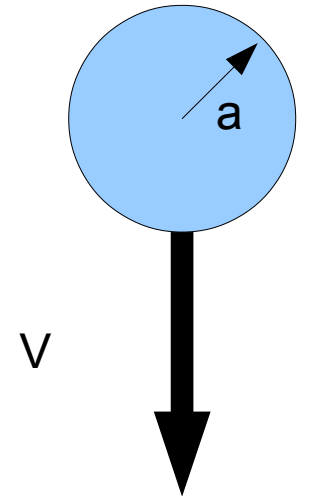


# Settling : final velocity of a droplet in the atmosphere

---

Sphere of radius  $a$  falling in the fluid have a final velocity of :

$$V_{stokes} = a^2 \frac{2g\Delta\rho}{9\eta} \quad \Delta\rho = \rho_{particle} - \rho_{fluid}$$



When the mean free path of the molecules becomes bigger than the size of the particle :  $\beta$

$$V_f = (1 + 1.26 K_N) \cdot V_{stokes}$$
$$K_N = \frac{\lambda}{a} \quad \lambda = \frac{k_B T}{\sqrt{2} \pi d^2} \cdot \frac{1}{P}$$

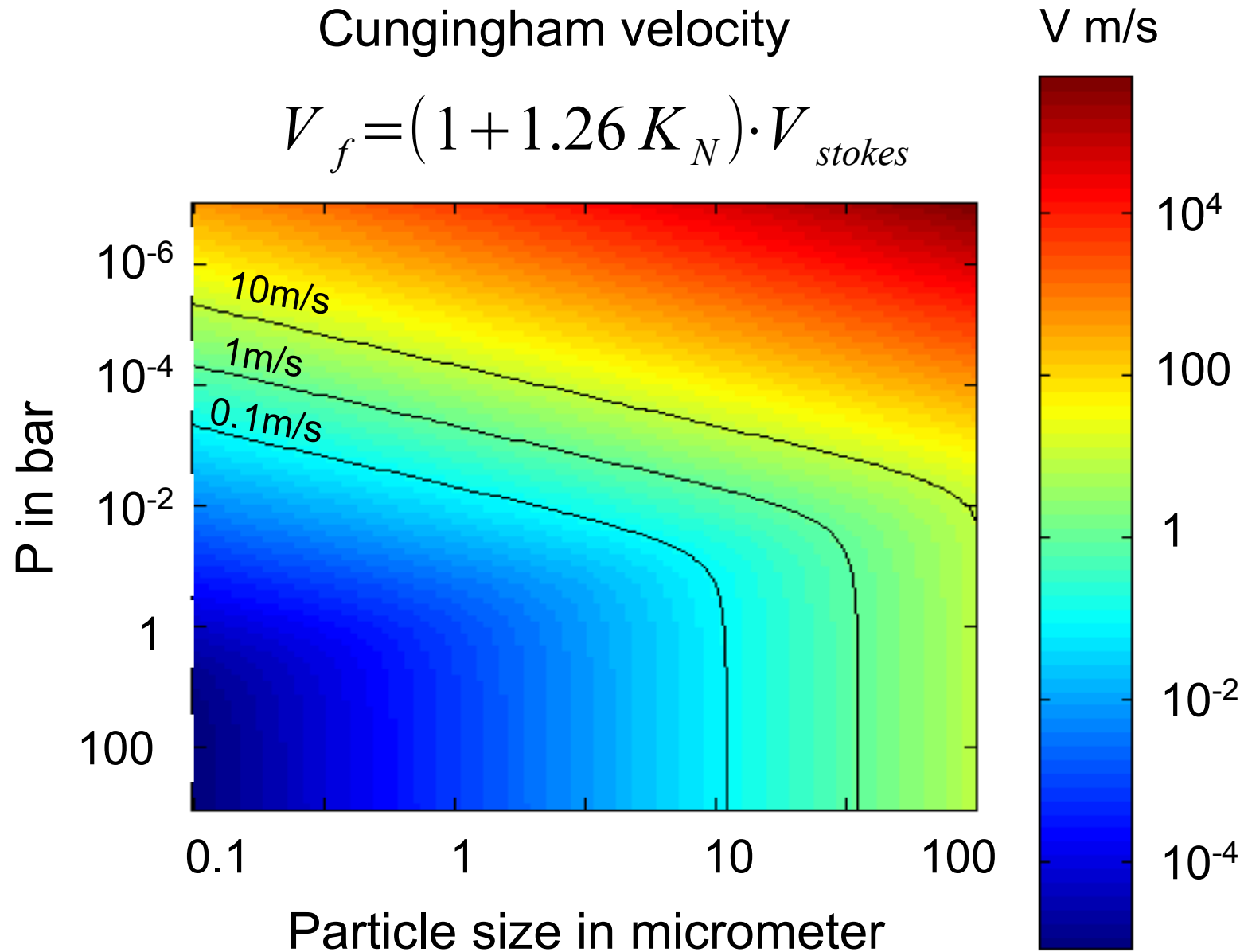
}  $\longrightarrow$   $V_f \propto \frac{a}{P}$

# Settling : final velocity of droplets in the atmosphere

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Cunningham velocity

$$V_f = (1 + 1.26 K_N) \cdot V_{stokes}$$



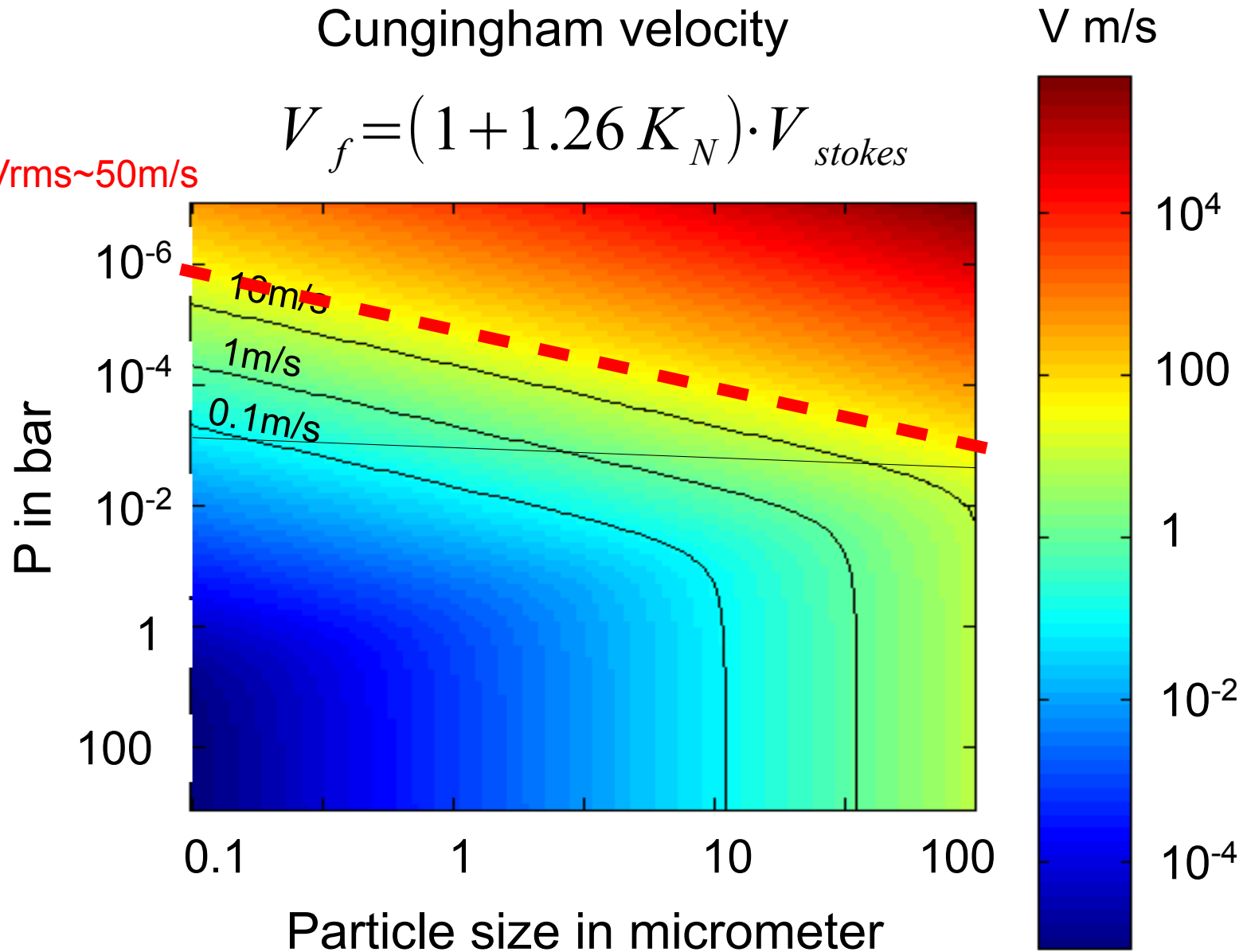


# Settling : final velocity of droplets in the atmosphere

Cunningham velocity

$$V_f = (1 + 1.26 K_N) \cdot V_{stokes}$$

$V_{rms} \sim 50 \text{ m/s}$



# Diffusion : equations

---


$$\frac{D\chi}{Dt} = \frac{1}{\rho} \frac{\partial \rho V_z \chi}{\partial z}$$

3D – model big  
simulations

...

A lot of approximations and simplifications

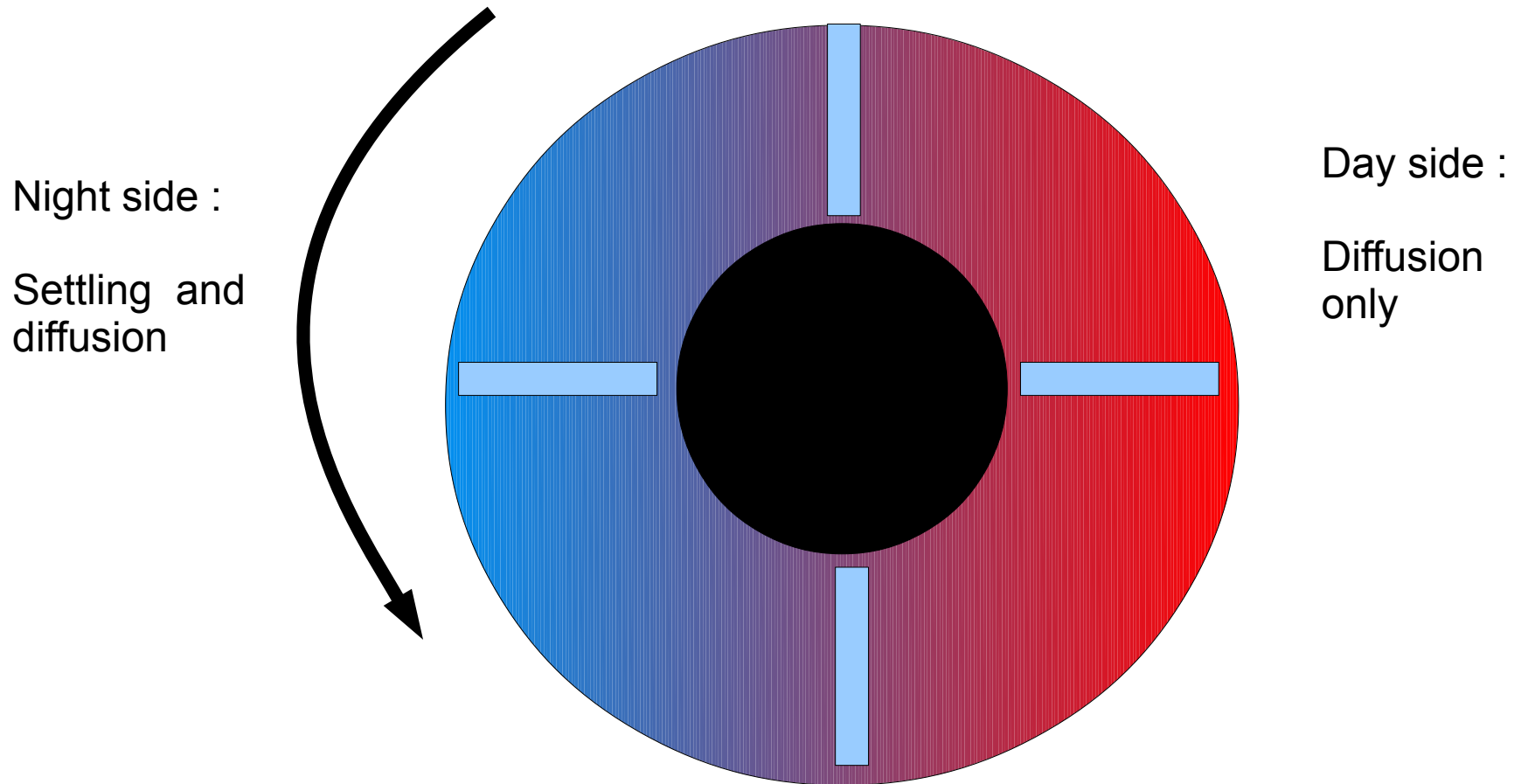
...


$$\frac{\partial \chi}{\partial t} - \frac{\partial}{\partial P} \left( K_P \frac{\partial \chi}{\partial P} \right) = \frac{\partial V_P \chi}{\partial P}$$

1D – highly idealized toy model

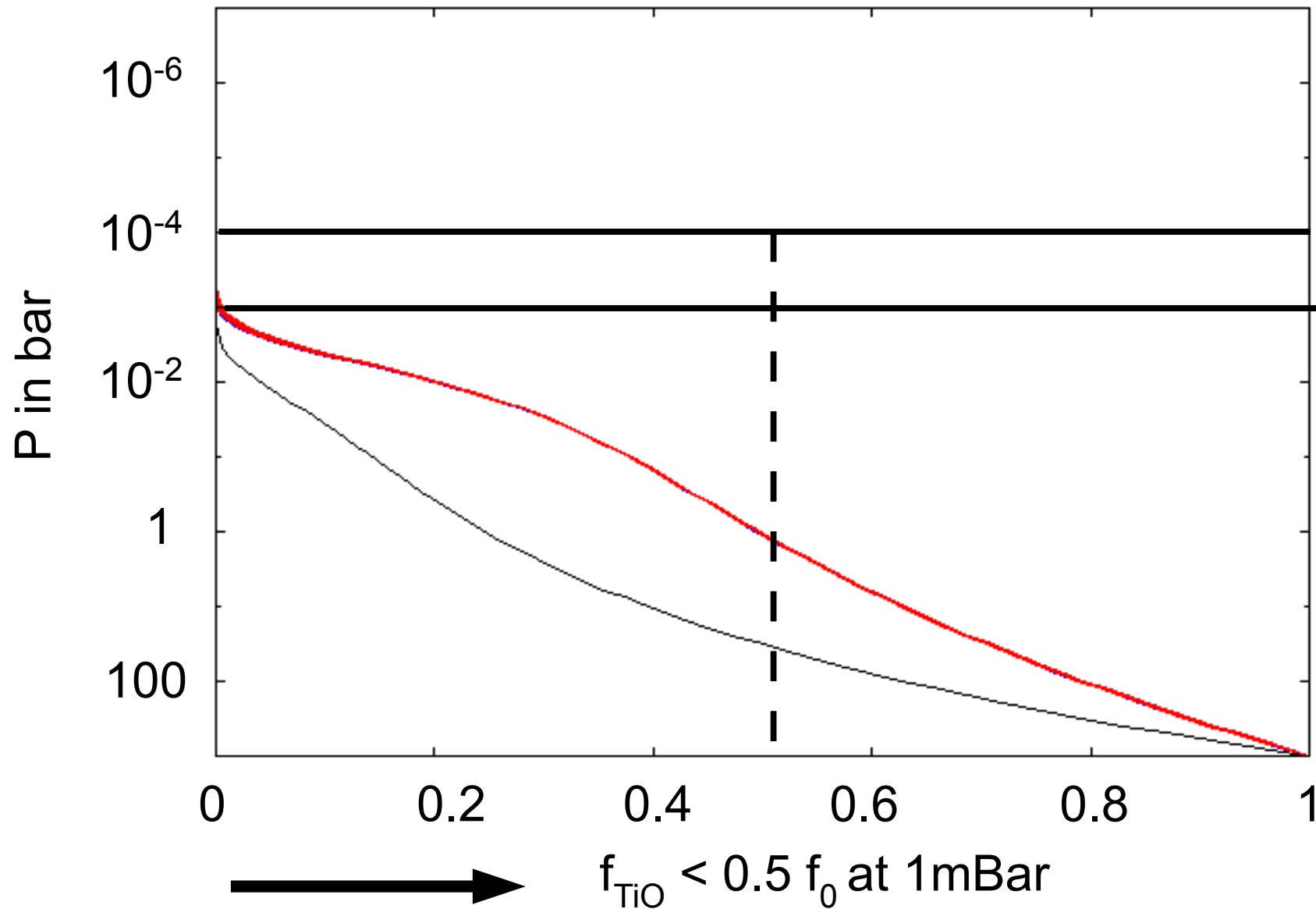
# Numerical model

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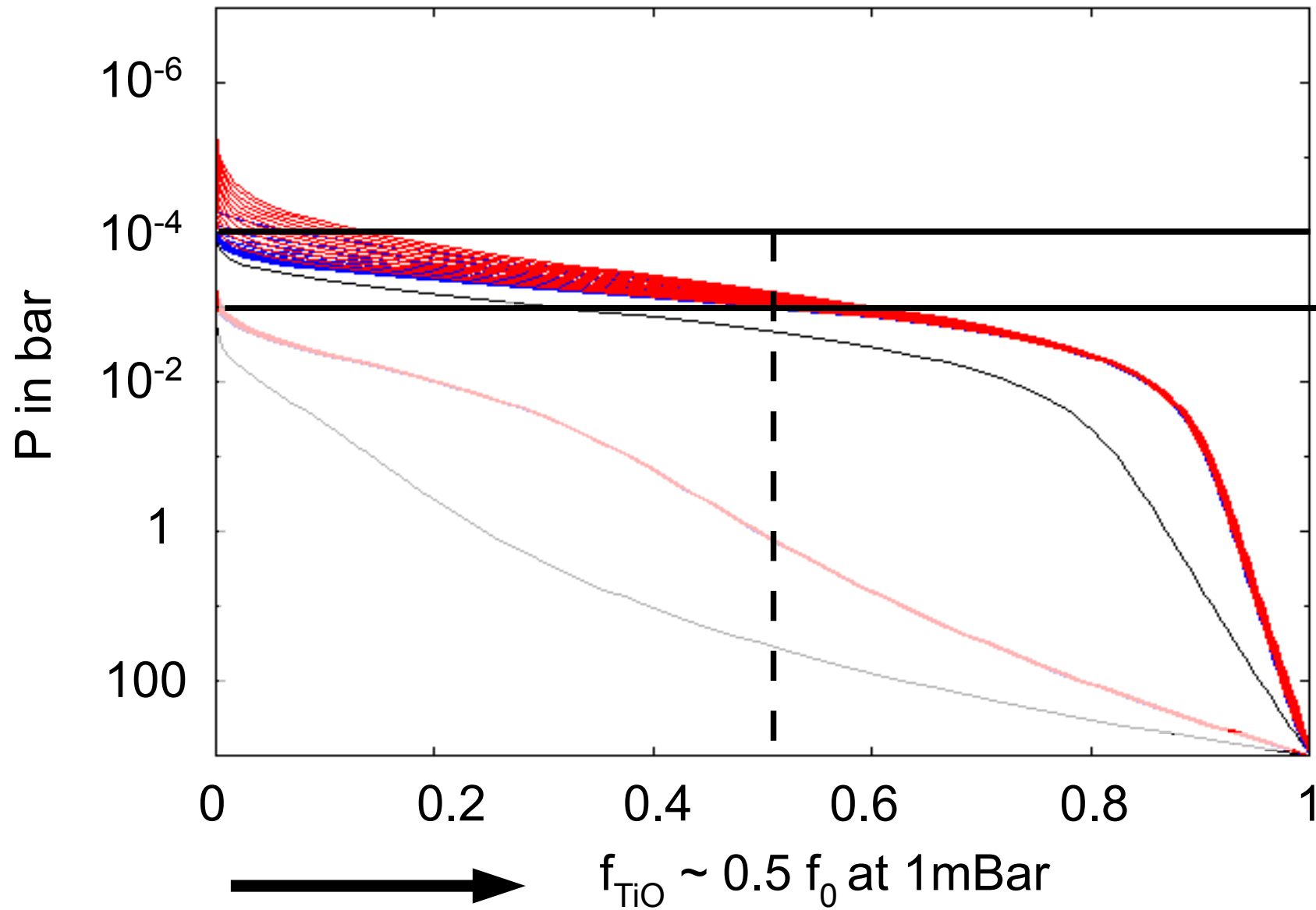
# Results : abundance profiles

$a=10\mu\text{m}$   $Kz=10^8\text{cm}^2/\text{s}$



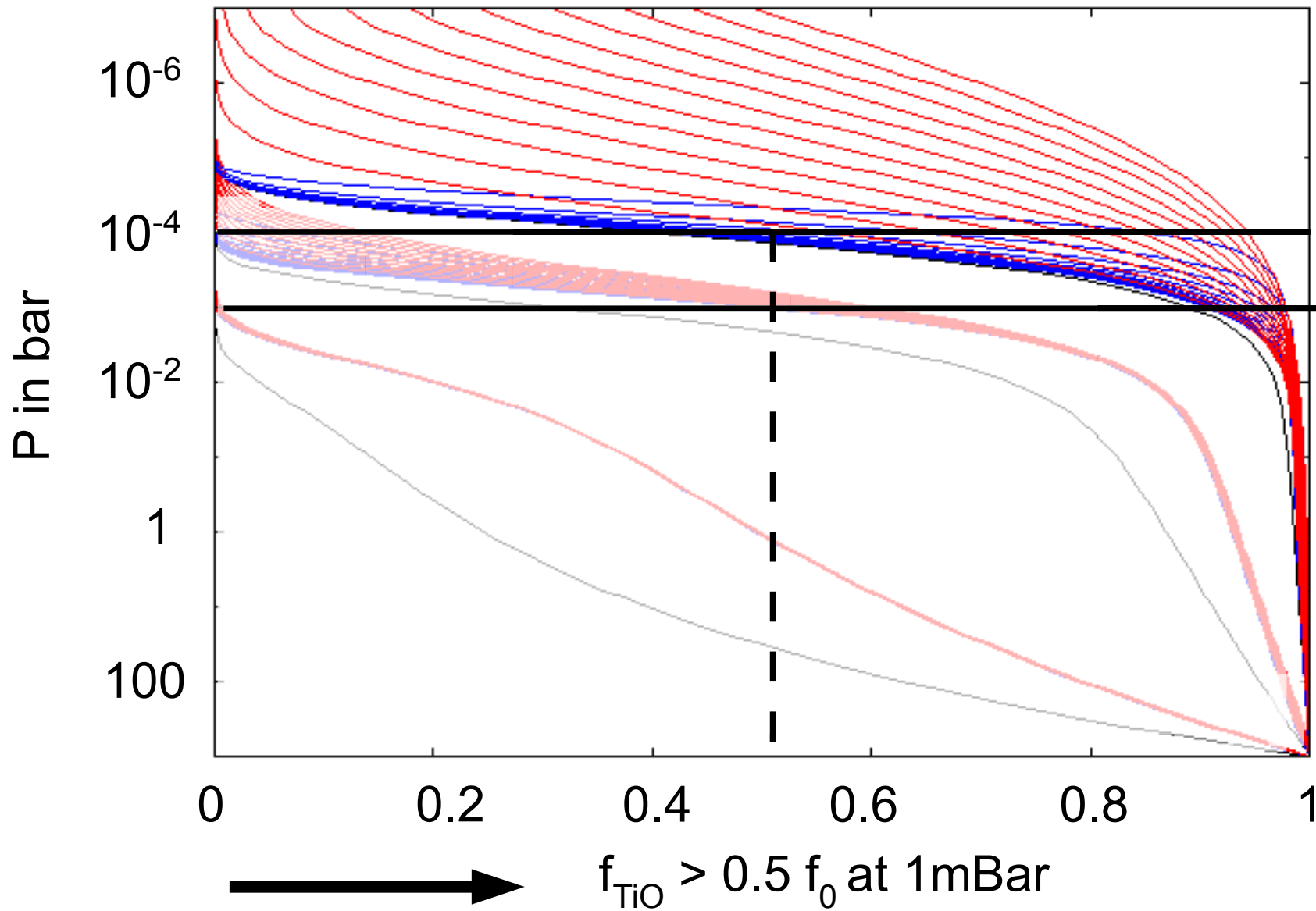
# Results : abundance profiles

$a=10\mu\text{m}$   $Kz=10^9\text{cm}^2/\text{s}$

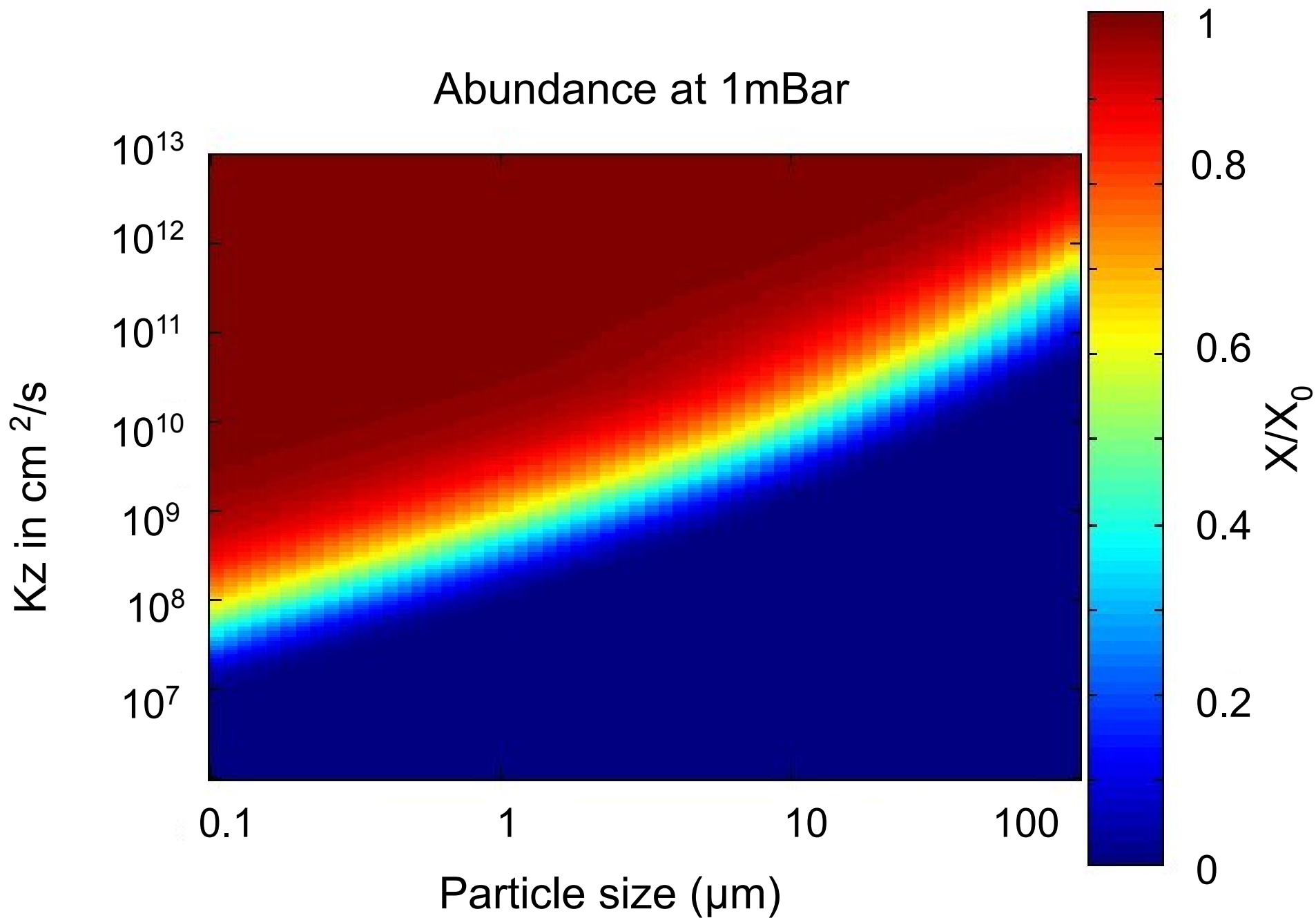


# Results : abundance profiles

$a=10\mu\text{m}$   $Kz=10^{10}\text{cm}^2/\text{s}$

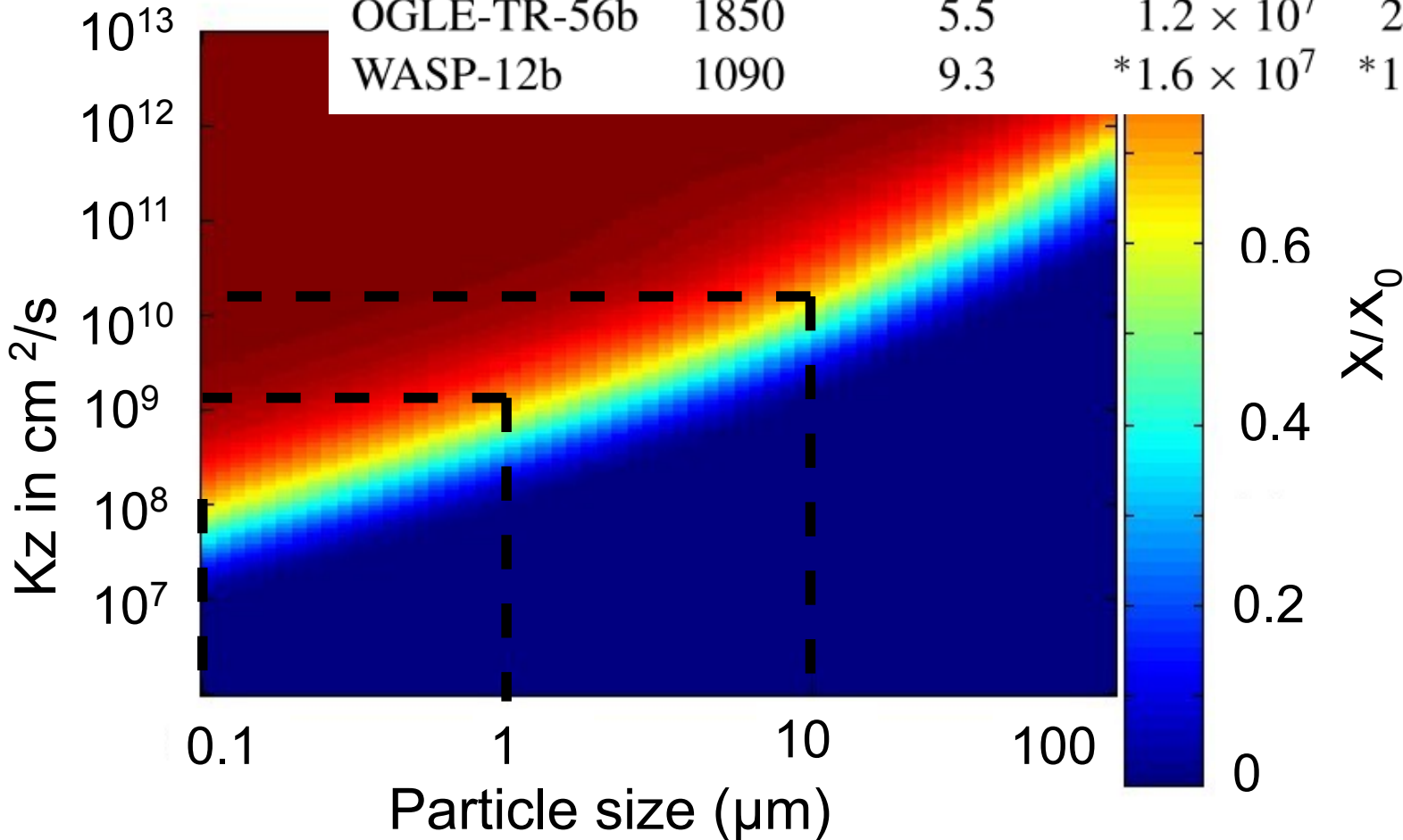


# Results: constrain on $K_z$



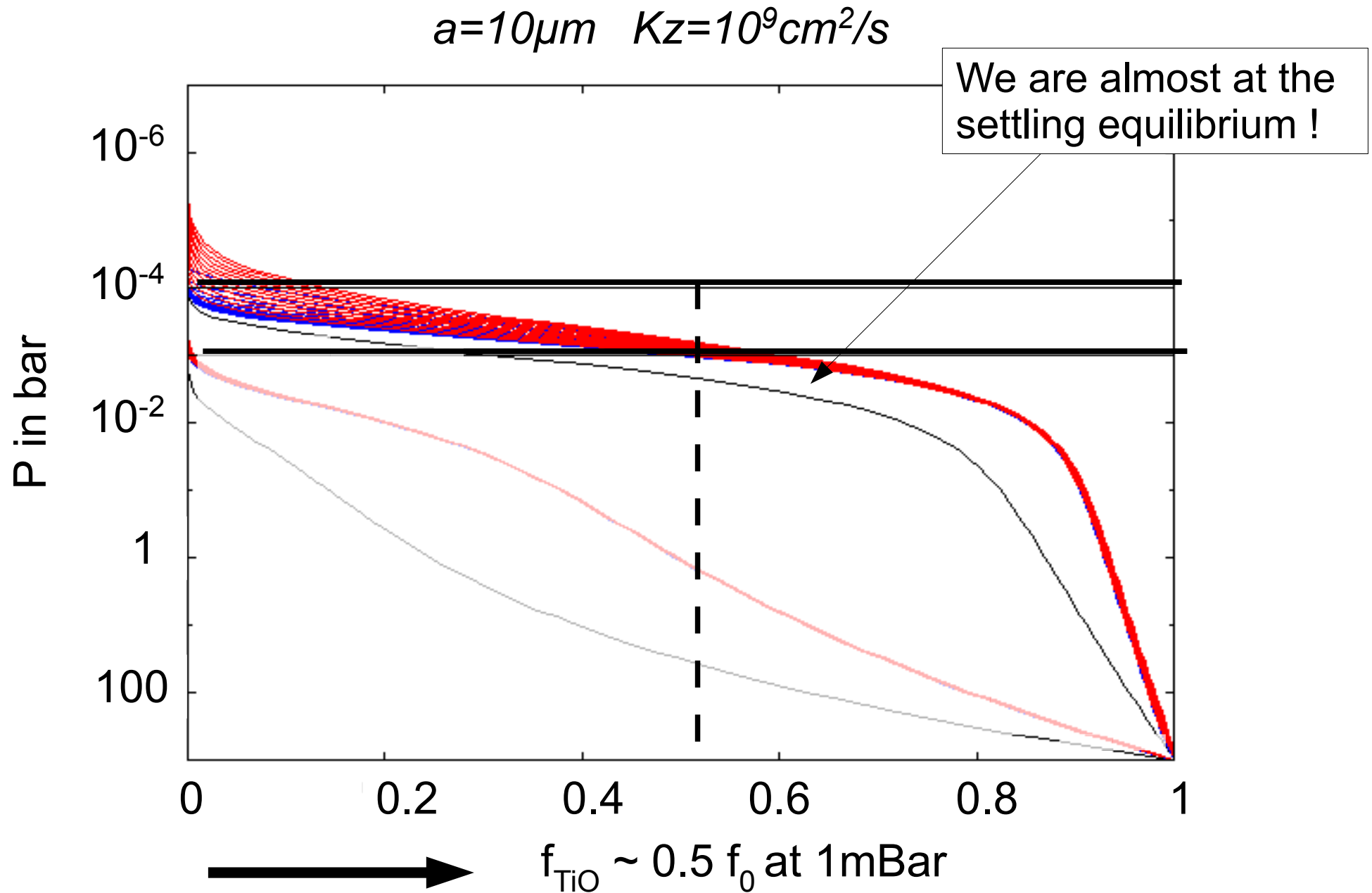
# Results: comparison with the vertical cold trap

Planet	$g$ ( $\text{cm s}^{-2}$ )	$F_*$ ( $\text{erg cm}^2 \text{s}^{-1}$ )	Required $K_{zz}$		
			$a = 0.1 \mu\text{m}$	$1 \mu\text{m}$	$10 \mu\text{m}$
HD 209458b	1000	1.0	$6.2 \times 10^8$	$6.2 \times 10^9$	$6.5 \times 10^{10}$
HD 149026b	1560	2.2	$2.4 \times 10^8$	$2.3 \times 10^9$	$2.6 \times 10^{10}$
TrES-4	721	2.4	$2.7 \times 10^8$	$2.7 \times 10^9$	$3.0 \times 10^{10}$
OGLE-TR-56b	1850	5.5	$1.2 \times 10^7$	$2.1 \times 10^7$	$8.7 \times 10^8$
WASP-12b	1090	9.3	* $1.6 \times 10^7$	* $1.6 \times 10^7$	* $1.6 \times 10^7$



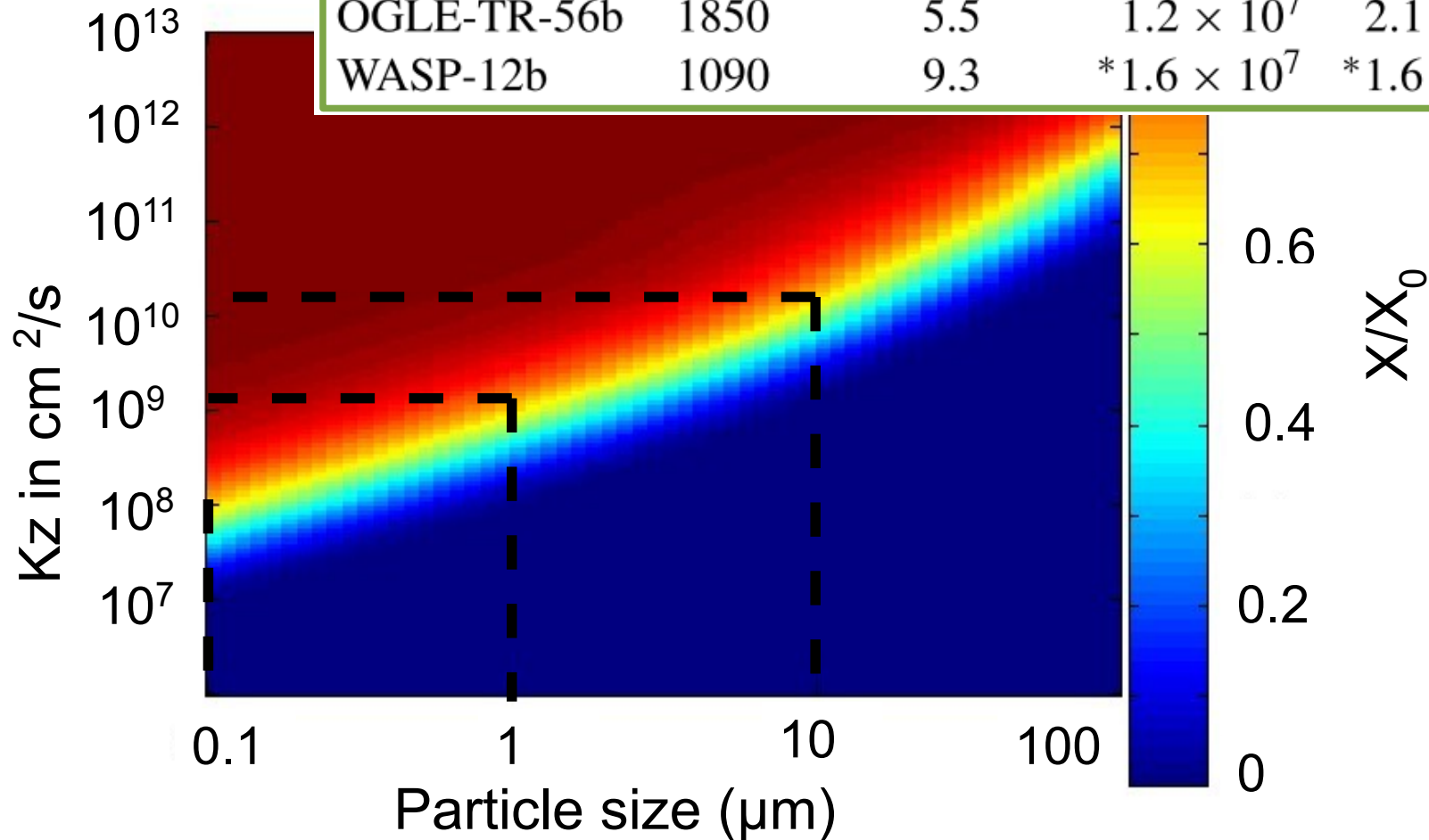


# Results : abundance profiles



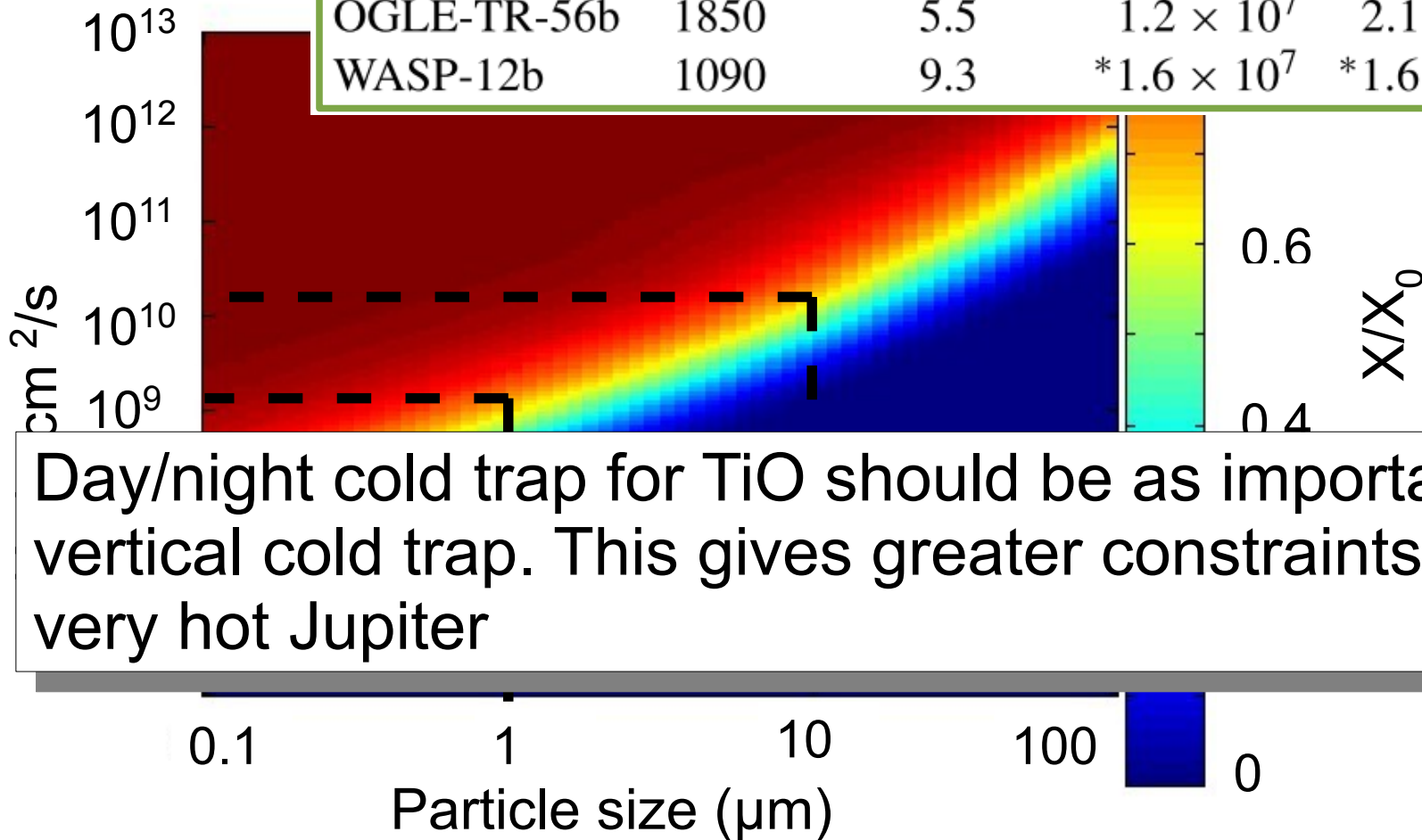
# Results: comparison with the vertical cold trap

Planet	$g$ ( $\text{cm s}^{-2}$ )	$F_*$ ( $\text{erg cm}^2 \text{s}^{-1}$ )	Required $K_{zz}$		
			$a = 0.1 \mu\text{m}$	$1 \mu\text{m}$	$10 \mu\text{m}$
HD 209458b	1000	1.0	$6.2 \times 10^8$	$6.2 \times 10^9$	$6.5 \times 10^{10}$
HD 149026b	1560	2.2	$2.4 \times 10^8$	$2.3 \times 10^9$	$2.6 \times 10^{10}$
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WASP-12b	1090	9.3	$*1.6 \times 10^7$	$*1.6 \times 10^7$	$*1.6 \times 10^7$



# Results: comparison with the vertical cold trap

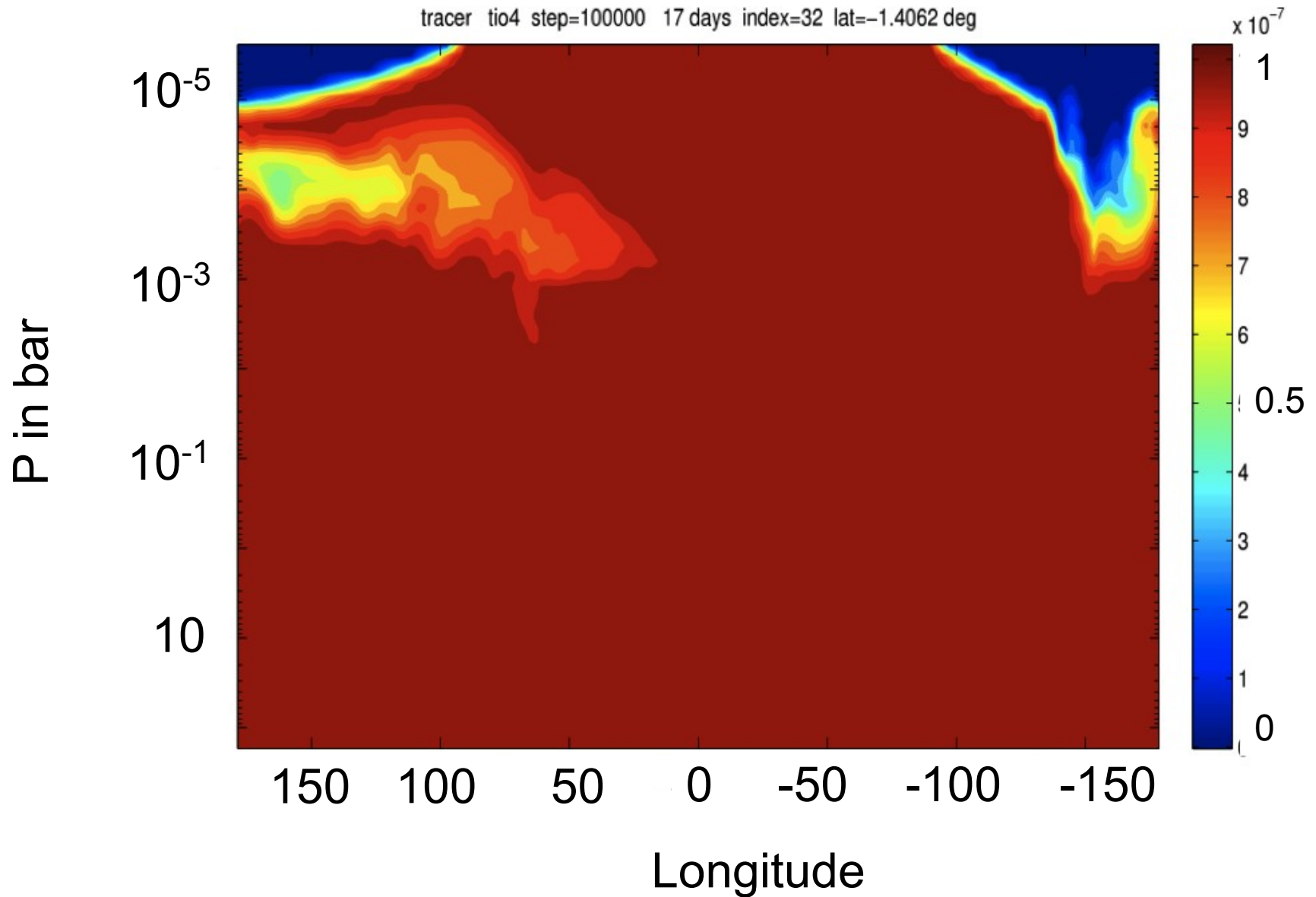
Planet	$g$ ( $\text{cm s}^{-2}$ )	$F_*$ ( $\text{erg cm}^2 \text{s}^{-1}$ )	Required $K_{zz}$		
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HD 209458b	1000	1.0	$6.2 \times 10^8$	$6.2 \times 10^9$	$6.5 \times 10^{10}$
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WASP-12b	1090	9.3	$*1.6 \times 10^7$	$*1.6 \times 10^7$	$*1.6 \times 10^7$



Day/night cold trap for TiO should be as important as vertical cold trap. This gives greater constraints on  $K_{zz}$  for very hot Jupiter

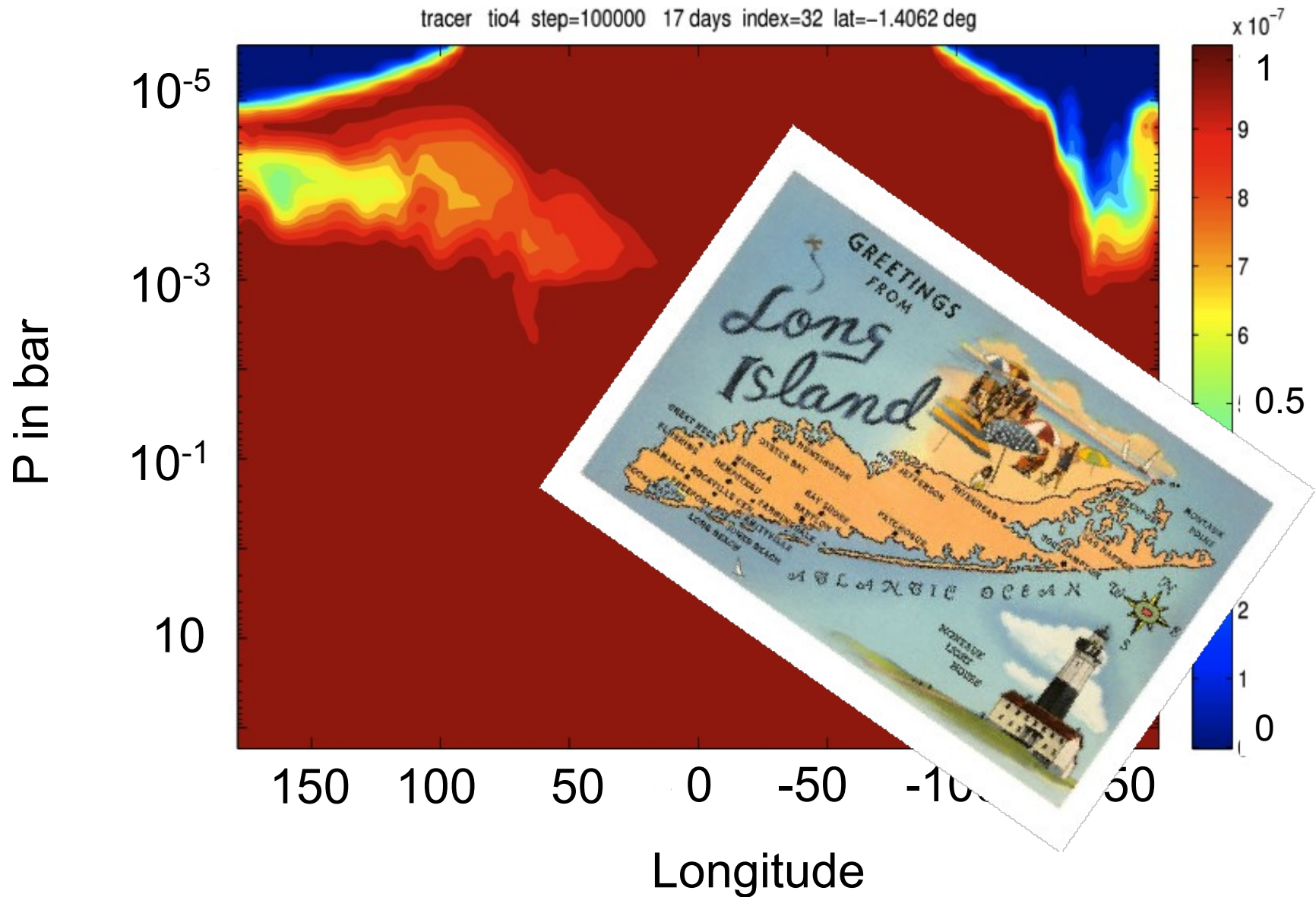
# Results : 3D simulation !

*Thank you Adam !*



# Results : 3D simulation !

*Thank you Adam !*



# Future work

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The presence or absence of TiO in hot Jupiter atmosphere depends strongly on the diffusivity of the flow



Measure the effective vertical mixing of the flow using tracers on a 3D simulation. Is the 1D mixing really diffusive ?



Different planets must have different conditions. We have to extend our work to hotter planet where the vertical cold trap is no more effective.



Run more simulations !